



Computational Homogenization and Multiscale Modeling

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Lecture 2: Contents

- Nested macro-micro computations (basis for FE^2)
 - Variationally consistent homogenization
 - Generalized macrohomogeneity (Hill-Mandel) condition
- Prolongation conditions - nonlinear SVE-problems
 - Boundary conditions versus macroscale strain or stress control – Overview (table)
 - Displacement Boundary Condition (nonlinear DBC-problem)
 - Traction Displacement Boundary Condition (nonlinear TBC-problem)
 - Weakly Periodic Boundary Condition (nonlinear PBC-problem) – generalization of classical strong displacement periodicity

Model-based homogenization

- Two separate scales *a priori* within each SVE. Local $\mathbf{u}(\bar{\mathbf{X}}; \mathbf{X}) \in \mathbb{U}$, global $\bar{\mathbf{u}}(\bar{\mathbf{X}}) \in \bar{\mathbb{U}}$

$$\mathbf{u}(\bar{\mathbf{X}}; \mathbf{X}) = \mathbf{u}^M(\bar{\mathbf{X}}; \mathbf{X}) + \mathbf{u}^s(\bar{\mathbf{X}}; \mathbf{X})$$

- Macroscale prolongation to Γ_{\square} . For given $\bar{\mathbf{X}} \in \Omega$:

$$\mathbf{u}^M(\mathbf{X}) = \mathbf{u}^M(\bar{\mathbf{H}}^{(0)} \stackrel{\text{def}}{=} \bar{\mathbf{u}}, \bar{\mathbf{H}}^{(1)} \stackrel{\text{def}}{=} \bar{\mathbf{H}}, \bar{\mathbf{H}}^{(2)}, \dots, \bar{\mathbf{H}}^{(K)}; \mathbf{X})$$

with "higher order" gradients (K = order of homogenization):

$$\bar{\mathbf{H}}^{(k)} \stackrel{\text{def}}{=} \bar{\mathbf{u}} \otimes \underbrace{\bar{\nabla} \otimes \bar{\nabla} \otimes \dots \otimes \bar{\nabla}}_k$$

- 1st order (conventional): $\mathbf{u}^M(\mathbf{X}) = \bar{\mathbf{u}} + \bar{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}]$
- 2nd order: $\mathbf{u}^M(\mathbf{X}) = \bar{\mathbf{u}} + \bar{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}] + \frac{1}{2}[\mathbf{X} - \bar{\mathbf{X}}] \cdot \bar{\mathbf{H}}^{(2)} \cdot [\mathbf{X} - \bar{\mathbf{X}}]$

Note: $\bar{\mathbf{H}}^{(2)}$ 3rd order tensor (such as curvature), e.g. KUZNETSOVA ET AL. 2002

Model-based homogenization

- Prolongation condition defines (implicitly) how the fluctuation depends on the macro-solution: $\mathbf{u}^s\{\mathbf{u}^M\}, \forall \mathbf{u}^M \in \mathbb{U}^M$

$$\Rightarrow \bar{\mathbf{u}}(\bar{\mathbf{X}}) \xrightarrow{\text{SVE}} \mathbf{u}^M(\bar{\mathbf{X}}, \mathbf{X}) \rightarrow \mathbf{u}^s\{\mathbf{u}^M(\bar{\mathbf{X}}, \mathbf{X})\} \rightarrow \mathbf{u}\{\bar{\mathbf{u}}(\bar{\mathbf{X}}), \mathbf{X}\}$$

- Homogenized problem:

$$y(\bar{\mathbf{X}}) \mapsto \langle y \rangle_{\square}(\bar{\mathbf{X}}), \quad \bar{\mathbf{X}} \in \Omega$$

Note: Global boundary data assumed to vary "slowly" $\Rightarrow y(\bar{\mathbf{X}}) \approx \bar{y}(\bar{\mathbf{X}})$ on Γ

- Variational problem: Find $\mathbf{u}^M \in \mathbb{U}^M$ s. t.

$$R(\mathbf{u}\{\mathbf{u}^M\}; \delta \mathbf{u}) \stackrel{\text{def}}{=} l(\delta \mathbf{u}) - a(\mathbf{u}\{\mathbf{u}^M\}; \delta \mathbf{u}) = 0 \quad \forall \delta \mathbf{u}^M \in \mathbb{U}^{M,0}.$$

where

$$a(\mathbf{u}; \delta \mathbf{u}) \stackrel{\text{def}}{=} \int_{\Omega} \langle \mathbf{P} : \delta \mathbf{H} \rangle_{\square} dV, \quad l(\delta \mathbf{u}) \stackrel{\text{def}}{=} \int_{\Omega} \langle \mathbf{f} \cdot \delta \mathbf{u} \rangle_{\square} dV + \int_{\Gamma_N} \bar{\mathbf{t}}_p \cdot \delta \bar{\mathbf{u}} dS$$

VMS vs. Homogenization

- Single-scale problem: Find $\mathbf{u} \in \mathbb{U}$ such that

$$R(\mathbf{u}; \delta \mathbf{u}) \stackrel{\text{def}}{=} l(\delta \mathbf{u}) - a(\mathbf{u}; \delta \mathbf{u}) = 0 \quad \forall \delta \mathbf{u} \in \mathbb{U}^0$$

- Variational Multiscale Method (HUGHES 1995) $(\mathbf{u}^M, \mathbf{u}^s) \in \mathbb{U}^M \times \mathbb{U}^s$

$$R(\mathbf{u}^M + \mathbf{u}^s; \delta \mathbf{u}^M) = 0 \quad \forall \delta \mathbf{u}^M \in \mathbb{U}^{M,0} \subset \mathbb{U}^0$$

$$R(\mathbf{u}^M + \mathbf{u}^s; \delta \mathbf{u}^s) = 0 \quad \forall \delta \mathbf{u}^s \rightarrow \mathbf{u}^s \{ \mathbf{u}^M \}$$

- Localized problem for $\mathbf{u}^s \Rightarrow$ We seek solutions $\tilde{\mathbf{u}}$

$$\leadsto \tilde{\mathbf{u}} \in \tilde{\mathbb{U}} \stackrel{\text{def}}{=} \{ \mathbb{U} \ni \mathbf{u} = \mathbf{u}^M + \mathbf{u}^s \{ \mathbf{u}^M \}, \mathbf{u}^M \in \mathbb{U}^M \}$$

- **Issue:** Due to approximations/restrictions

$$R(\tilde{\mathbf{u}}; \delta \mathbf{u}) \neq 0$$

In general not even for $\delta \mathbf{u} = \delta \mathbf{u}^M + \mathbf{u}^s \{ \delta \mathbf{u}^M \}$

Two possible choices for weighted residuals

I Classical assumption (in VMS *and* Homogenization)

$$R(\mathbf{u}^M + \mathbf{u}^s; \delta \mathbf{u}^M) = 0 \quad \forall \delta \mathbf{u}^M \in \mathbb{U}^{M,0}$$

- Straightforward choice (presuming globally exact $\mathbf{u}^s \{\bullet\}$, i.e. $\mathbb{U}^s = \mathbb{U} \ominus \mathbb{U}^M$)
- $\delta \mathbf{u}^M$ "homogenizer" \rightsquigarrow Balance of homogenized fluxes/stresses

II Generalized (non-linear) Galerkin-type

$$R(\underbrace{\mathbf{u}^M + \mathbf{u}^s}_{\tilde{\mathbf{u}}}; \underbrace{\delta \mathbf{u}^M + (\mathbf{u}^s)' \{\mathbf{u}^M; \delta \mathbf{u}^M\}}_{\delta \tilde{\mathbf{u}}}) = 0 \quad \forall \delta \mathbf{u}^M \in \mathbb{U}^{M,0}$$

Test space: tangent space to approximation space

$$\delta \tilde{\mathbf{u}} \in \tilde{\mathbb{U}}^0(\mathbf{u}^M) \stackrel{\text{def}}{=} \left\{ \mathbb{U} \ni \delta \mathbf{u} = \delta \mathbf{u}^M + \underbrace{(\mathbf{u}^s)' \{\mathbf{u}^M; \delta \mathbf{u}^M\}}_{\text{sensitivity}}, \delta \mathbf{u}^M \in \mathbb{U}^M \right\}$$

- Variation of energy (if \exists)
- Preserves symmetry for $R'(\mathbf{u}; \delta \mathbf{u}, \delta \mathbf{u})$

Variationally Consistent Macrohomogeneity Condition (VCMC)

- By construction of local problems, i.e. by defining $\mathbf{u}^s\{\bullet\}$

$$\underbrace{R(\mathbf{u}^M + \mathbf{u}^s; \delta\mathbf{u}^M)}_{\text{I}} - \underbrace{R(\mathbf{u}^M + \mathbf{u}^s; \delta\mathbf{u}^M + (\mathbf{u}^s)' \{\mathbf{u}^M; \delta\mathbf{u}^M\})}_{\text{II}}$$

$$= \underbrace{R(\mathbf{u}^M + \mathbf{u}^s; (\mathbf{u}^s)' \{\mathbf{u}^M; \delta\mathbf{u}^M\})}_{\text{Fluctuation residual}} = 0 \quad \forall \delta\mathbf{u}^M \in \mathbb{U}^{M,0}$$

- Solve **I** and get **II** "for free"
 - Energy equivalence (Hill-Mandel condition)
 - Symmetric VMS-problem

Macrohomogeneity condition

- "Local" version of VCMC in terms of SVE-residual

$$R(\mathbf{u}; \delta \mathbf{u}^s) \stackrel{\text{def}}{=} \int_{\Omega} R_{\square}(\mathbf{u}; \delta \mathbf{u}^s) dV$$

$$\begin{aligned} R_{\square}(\mathbf{u}; \delta \mathbf{u}^s) &\stackrel{\text{def}}{=} \langle \mathbf{f} \cdot \delta \mathbf{u}^s \rangle_{\square} - \langle \mathbf{P} : \delta \mathbf{H}^s \rangle_{\square} = \frac{1}{|\Omega_{\square}|} \int_{\Omega_{\square}} [\mathbf{f} \cdot \delta \mathbf{u}^s - \mathbf{P} : \delta \mathbf{H}^s] dV \\ &= -\frac{1}{|\Omega_{\square}|} \int_{\Gamma_{\square}} \mathbf{t} \cdot \delta \mathbf{u}^s dS \end{aligned}$$

- Sufficient condition for VCMC

$$R_{\square}(\mathbf{u}\{u^M\}; (\mathbf{u}^s)' \{u^M; \delta u^M\}) = 0, \quad \forall \delta u^M \in \mathbb{U}^{M,0}.$$

- **Note:** No need to solve for sensitivity fields

VCMC - 1st order homogenization

- VCMC \equiv HMC (classical Hill-Mandel macrohomogeneity condition)

$$\mathbf{u}^M(\mathbf{X}) = \bar{\mathbf{u}} + \bar{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}]$$

$$\Rightarrow \delta \mathbf{u}^M(\mathbf{X}) = \delta \bar{\mathbf{u}} + \delta \bar{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}], \quad \delta \mathbf{H}^M(\mathbf{X}) = \delta \bar{\mathbf{H}}$$

$$\Rightarrow \langle \mathbf{P} : \delta \mathbf{H} \rangle_{\square} = \langle \mathbf{P} \rangle_{\square} : \langle \delta \mathbf{H} \rangle_{\square} = \bar{\mathbf{P}} : \delta \bar{\mathbf{H}}, \quad \forall \delta \mathbf{u} \in \mathbb{U}_{\square}$$

- **Note:** HMC valid only if the condition $\langle \delta \mathbf{H} \rangle_{\square} = \delta \bar{\mathbf{H}}$ is satisfied:
 - automatically by the choice of test functions
 - imposed as part of the SVE-problem
- **Note:** Higher order homogenization "stresses" and "strains" on macro-level will appear in case of higher order homogenization, e.g. GEERS ET AL.

Explicit homogenization results

- Homogenized variational forms

$$a(\mathbf{u}\{\mathbf{u}^M\}; \delta\mathbf{u}^M) = \int_{\Omega} \langle \mathbf{P} : \delta\mathbf{H}^M \rangle_{\square} d\bar{V},$$

$$l(\delta\mathbf{u}^M) = \int_{\Omega} \langle \mathbf{f} \cdot \delta\mathbf{u}^M \rangle_{\square} d\bar{V} + \int_{\Gamma_N} \bar{\mathbf{t}}_p \cdot \delta\bar{\mathbf{u}} d\bar{S}$$

$$\text{SVE-forms: } \langle \mathbf{P} : \delta\mathbf{H} \rangle_{\square} = \bar{\mathbf{P}}\{\bar{\mathbf{H}}\} : \delta\bar{\mathbf{H}}, \quad \langle \mathbf{f} \cdot \delta\mathbf{u}^M \rangle_{\square} = \bar{\mathbf{f}} \cdot \delta\bar{\mathbf{u}} + \underbrace{\bar{\mathbf{f}}^{(2)} \cdot \delta\bar{\mathbf{H}}}_{\text{ignored}}$$

- Macroscale problem: Find $\bar{\mathbf{u}} \in \bar{\mathbb{U}}$ s. t.

$$\bar{R}\{\bar{\mathbf{u}}; \delta\bar{\mathbf{u}}\} \stackrel{\text{def}}{=} \bar{l}\{\delta\bar{\mathbf{u}}\} - \bar{a}\{\bar{\mathbf{u}}; \delta\bar{\mathbf{u}}\} = 0 \quad \forall \delta\bar{\mathbf{u}} \in \bar{\mathbb{U}}^0.$$

$$\bar{a}\{\bar{\mathbf{u}}; \delta\bar{\mathbf{u}}\} \stackrel{\text{def}}{=} \int_{\Omega} \bar{\mathbf{P}}\{\bar{\mathbf{H}}\} : \delta\bar{\mathbf{H}} d\bar{V}, \quad \bar{l}\{\delta\bar{\mathbf{u}}\} \stackrel{\text{def}}{=} \int_{\Omega} \bar{\mathbf{f}} \cdot \delta\bar{\mathbf{u}} d\bar{V} + \int_{\Gamma_N} \bar{\mathbf{t}}_p \cdot \delta\bar{\mathbf{u}} d\bar{S}$$

- **Note:** \bar{a} may represent higher order continuum model

Nested FE²-algorithm

- Solution via Newton iterations based on the tangent form

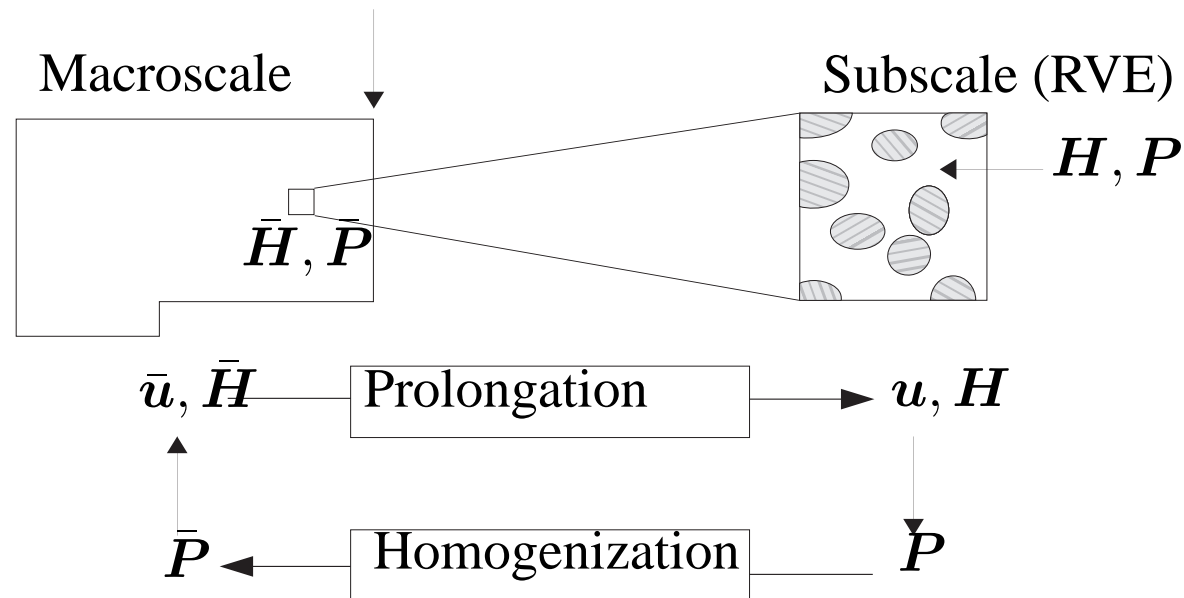
$$\begin{aligned} \bar{a}'\{\bar{\mathbf{u}}; \delta\bar{\mathbf{u}}, \Delta\bar{\mathbf{u}}\} &\stackrel{\text{def}}{=} \int_{\Omega} a'_{\square}(\mathbf{u}\{\mathbf{u}^{\text{M}}\}; \delta\mathbf{u}^{\text{M}}, \Delta\mathbf{u}^{\text{M}} + (\mathbf{u}^{\text{s}})'\{\mathbf{u}^{\text{M}}; \Delta\mathbf{u}^{\text{M}}\}) \, d\bar{V} \\ &= \int_{\Omega} \delta\bar{\mathbf{H}} : \bar{\mathbf{L}}_{\text{AT}} : \Delta\bar{\mathbf{H}} \, d\bar{V} \end{aligned}$$

Macroscale Algorithmic Tangent (AT) stiffness $\bar{\mathbf{L}}_{\text{AT}}$ defined from

$$d\bar{\mathbf{P}}\{\bar{\mathbf{H}}\} = \bar{\mathbf{L}}_{\text{AT}}\{\bar{\mathbf{H}}\} : d\bar{\mathbf{H}}$$

- $\bar{\mathbf{L}}_{\text{AT}}$ obtained from linearized SVE-problem (sensitivity problem)
- Sensitivity problem formulated for each choice of prolongation condition

Nested FE²-algorithm



1. Assume (nonequilibrium) macroscale stress \bar{P} , in the macroscale iteration, corresponding to given \bar{u} (and \bar{H}).
2. **Prolongation-homogenization:** For given $\bar{H} = \bar{H}(\bar{X}_i)$ in each macroscale (Gauss) quadrature point \bar{X}_i , solve the nonlinear RVE-problem (with chosen prolongation conditions) for subscale $u(\bar{X}_i, X)$ and homogenized (macroscale) stress $\bar{P}(\bar{X}_i)$.
3. Check macroscale residual. If convergence, $|\bar{R}| < TOL$, then stop, else compute a new (updated) \bar{u} (while using \bar{L}_{AT}) and return to 1.

DBC – Macroscale strain control

- Appropriate subscale spaces

$$\begin{aligned} \mathbb{U}_{\square}^{\text{D}} &= \{ \mathbf{u} \in \mathbb{U}_{\square} \mid \exists \hat{\mathbf{H}} \in \mathbb{R}^{3 \times 3} \text{ s.t. } \mathbf{u} = \hat{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}] \text{ on } \Gamma_{\square} \} \\ \mathbb{U}_{\square}^{\text{D},0} &= \{ \mathbf{u} \in \mathbb{U}_{\square} \mid \mathbf{u} = \mathbf{0} \text{ on } \Gamma_{\square} \} \end{aligned}$$

- SVE-problem: For given value of $\bar{\mathbf{H}} \in \mathbb{R}^{3 \times 3}$, find $\mathbf{u}^{\text{s}} \in \mathbb{U}_{\square}^{\text{D},0}$ that solves

$$a_{\square}(\mathbf{u}^{\text{M}}(\bar{\mathbf{H}}) + \mathbf{u}^{\text{s}}; \delta \mathbf{u}^{\text{s}}) = 0 \quad \forall \delta \mathbf{u}^{\text{s}} \in \mathbb{U}_{\square}^{\text{D},0}.$$

with $\mathbf{u}^{\text{M}}(\bar{\mathbf{H}}) = \bar{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}]$.

Solved by Newton iterations in standard fashion

- Postprocessing: $\bar{\mathbf{P}} = \langle \mathbf{P} \rangle_{\square}$

DBC– Macroscale strain control: AT stiffness

- Macroscale Algorithmic Tangent (AT) stiffness tensor: Compute sensitivity in \mathbf{u}^s for change in $\bar{\mathbf{H}}$:

$$d\bar{\mathbf{H}} \rightarrow d\mathbf{u} = d\mathbf{u}^M + d\mathbf{u}^s = d\mathbf{u}^M + (\mathbf{u}^s)' \{ \mathbf{u}^M; d\mathbf{u}^M \}$$

- *Ansatz*: $d\mathbf{u}^s = \sum_{i,j} \hat{\mathbf{u}}^{s(ij)} d\bar{H}_{ij}$
- "Unit fluctuation fields" (sensitivity fields) $\hat{\mathbf{u}}^{s(kl)}$ for $k, l = 1, 2, \dots, NDIM$

$$(a_{\square})'(\bullet; \delta\mathbf{u}^s, \hat{\mathbf{u}}^{s(kl)}) = -(a_{\square})'(\bullet; \delta\mathbf{u}^s, \hat{\mathbf{u}}^{M(kl)}) \quad \forall \delta\mathbf{u}^s \in \mathbb{U}_{\square}^{\mathbf{D},s}$$

- Primal approach for computing AT-tensor

$$\bar{\mathbf{L}}_{\text{AT}} = \langle \mathbf{L} \rangle_{\square} + \sum_{k,l=1}^{NDIM} \left\langle \mathbf{L} : \hat{\mathbf{H}}^{s(kl)} \right\rangle_{\square} \otimes \mathbf{e}_k \otimes \mathbf{e}_l$$

- Dual approach for computing AT-tensor: Useful method in conjunction with goal-oriented error computation for subscale FE-discretization, LARSSON AND RUNESSON 2007, "power of duality"!

TBC – Macroscale strain control

- Appropriate subscale spaces

$$\begin{aligned}\mathbb{U}_{\square}^{\text{N}} &= \{\mathbf{u} \in \mathbb{U}_{\square} \text{ s.t. } \mathbf{u}(\bar{\mathbf{X}}) = \mathbf{0}\} \\ \mathbb{U}_{\square}^{\text{N,s}} &= \{\mathbf{u} \in \mathbb{U}_{\square} : \langle \mathbf{H} \rangle_{\square} = \mathbf{0}\}\end{aligned}$$

- SVE-problem: For given value of $\bar{\mathbf{H}} \in \mathbb{R}^{3 \times 3}$, find $\mathbf{u} \in \mathbb{U}_{\square}^{\text{N}}$ and $\bar{\mathbf{P}} \in \mathbb{R}^{3 \times 3}$ that solve

$$\begin{aligned}a_{\square}(\mathbf{u}; \delta \mathbf{u}) - c_{\square}^{(\text{H})}(\delta \mathbf{u}; \bar{\mathbf{P}}) &= 0 \quad \forall \delta \mathbf{u} \in \mathbb{U}_{\square}^{\text{N}} \\ c_{\square}^{(\text{H})}(\mathbf{u}; \delta \bar{\mathbf{P}}) &= \delta \bar{\mathbf{P}} : \bar{\mathbf{H}} \quad \forall \delta \bar{\mathbf{P}} \in \mathbb{R}^{3 \times 3}\end{aligned}$$

where

$$c_{\square}^{(\text{H})}(\mathbf{u}; \bar{\mathbf{P}}) \stackrel{\text{def}}{=} \langle \mathbf{H} \rangle_{\square} : \bar{\mathbf{P}}$$

- **Note:** Term $c_{\square}^{(\text{H})}(\delta \mathbf{u}; \bar{\mathbf{P}})$ follows from model assumption

$$\mathbf{t} = \mathbf{P} \cdot \mathbf{N} = \bar{\mathbf{P}} \cdot \mathbf{N} \quad \text{in } \Gamma_{\square}$$

- No postprocessing of $\bar{\mathbf{P}} = \langle \mathbf{P} \rangle_{\square}$, "built-in" property in variational format

TBC – Macroscale strain control: AT stiffness

- Macroscale Algorithmic Tangent (AT) stiffness tensor: Compute sensitivity in \mathbf{u} and $\bar{\mathbf{P}}$ for change in $\bar{\mathbf{H}}$:

$$d\bar{\mathbf{H}} \rightarrow d\mathbf{u}, d\bar{\mathbf{P}}$$

- Ansatz: $d\mathbf{u}^s = \sum_{i,j} \hat{\mathbf{u}}^{s(ij)} d\bar{H}_{ij}$, $d\bar{\mathbf{P}} = \sum_{i,j} \hat{\bar{\mathbf{P}}}^{(ij)} d\bar{H}_{ij}$
- "Unit fluctuation fields" (sensitivity fields) $\hat{\mathbf{u}}^{s(kl)}$ and $\hat{\bar{\mathbf{P}}}^{(ij)}$ for $k, l = 1, 2, \dots, NDIM$

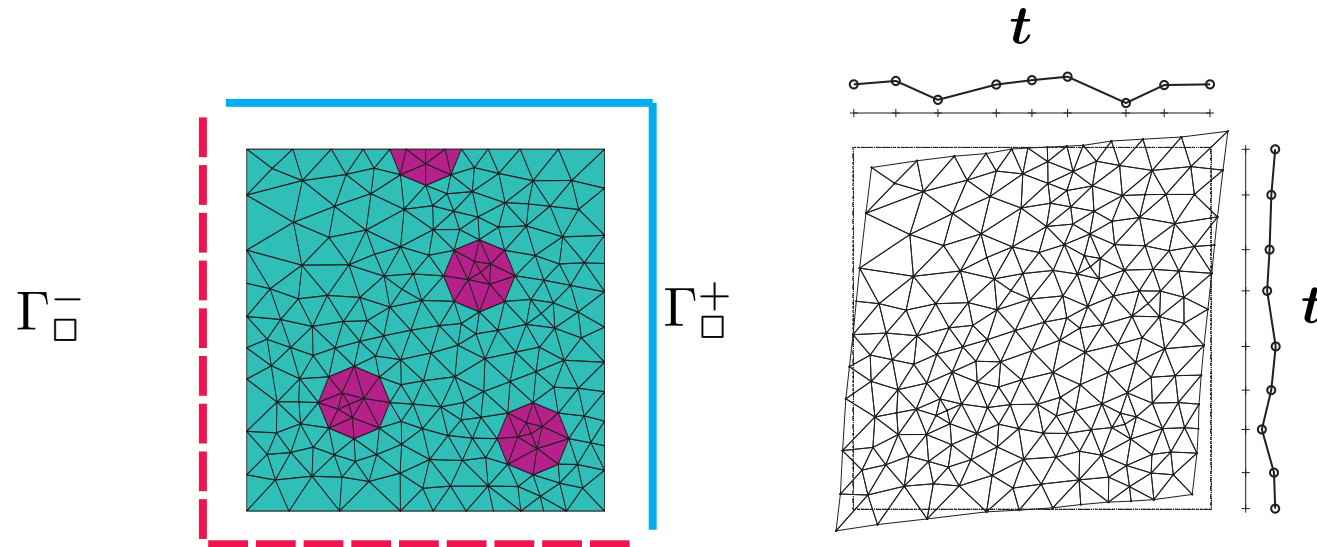
$$\begin{aligned} (a_{\square})'(\bullet; \delta\mathbf{u}, \hat{\mathbf{u}}^{(kl)}) - c_{\square}^{(H)}(\delta\mathbf{u}, \hat{\bar{\mathbf{P}}}^{(kl)}) &= 0 \quad \forall \delta\mathbf{u} \in \mathbb{U}_{\square}^N \\ c_{\square}^{(H)}(\hat{\mathbf{u}}^{(kl)}, \delta\bar{\mathbf{P}}) &= (\delta\bar{\mathbf{P}})_{kl} \quad \forall \delta\bar{\mathbf{P}} \in \mathbb{R}^{3 \times 3} \end{aligned}$$

- AT-tensor

$$\bar{\mathbf{L}}_{\text{AT}} = \sum_{k,l=1}^{NDIM} \hat{\bar{\mathbf{P}}}^{(kl)} \otimes \mathbf{e}_k \otimes \mathbf{e}_l$$

- **Note:** MIEHE AND KOCH 2002 in the FE-discrete setting, LARSSON AND RUNESSON, 2007 in the continuous setting

WBC (weak periodicity) – Macroscale strain control



- Periodic boundary conditions using "mirroring" ($\mathbf{X}^- : \Gamma_\square^+ \rightarrow \Gamma_\square^-$), $\Gamma_\square = \Gamma_\square^+ \cup \Gamma_\square^-$
- Periodic displacements and anti-periodic tractions

$$[[\mathbf{u}]] \stackrel{\text{def}}{=} \mathbf{u}(\mathbf{X}) - \mathbf{u}(\mathbf{X}^-(\mathbf{X})) = \bar{\mathbf{H}} \cdot [[\mathbf{X}]] \quad \text{on } \Gamma_\square^+$$

$$\mathbf{t} \stackrel{\text{def}}{=} \mathbf{t}(\mathbf{X}) = -\mathbf{t}(\mathbf{X}^-(\mathbf{X})) \quad \text{on } \Gamma_\square^+$$

WBC cont'd

- Appropriate subscale spaces

$$\mathbb{U}_\square = \mathbb{U}_\square^N = \{\mathbf{u} \in \mathbb{U}_\square \text{ s.t. } \mathbf{u}(\bar{\mathbf{X}}) = \mathbf{0}\}$$

$$\mathbb{T}_\square = \{\mathbf{t} \text{ "suff. regular"}\} \quad \textbf{Note: Built-in self-equilibration}$$

- Weak micro-periodicity formulation for SVE: For given value of $\bar{\mathbf{H}}$, find $\mathbf{u}, \mathbf{t} \in \mathbb{U}_\square \times \mathbb{T}_\square$ s.t.

$$\begin{aligned} a_\square(\mathbf{u}; \delta\mathbf{u}) - d_\square(\mathbf{t}, \delta\mathbf{u}) &= 0 & \forall \delta\mathbf{u} \in \mathbb{U}_\square \\ -d_\square(\delta\mathbf{t}, \mathbf{u}) &= -d_\square(\delta\mathbf{t}, \bar{\mathbf{H}} \cdot [\mathbf{X} - \bar{\mathbf{X}}]) & \forall \delta\mathbf{t} \in \mathbb{T}_\square \end{aligned}$$

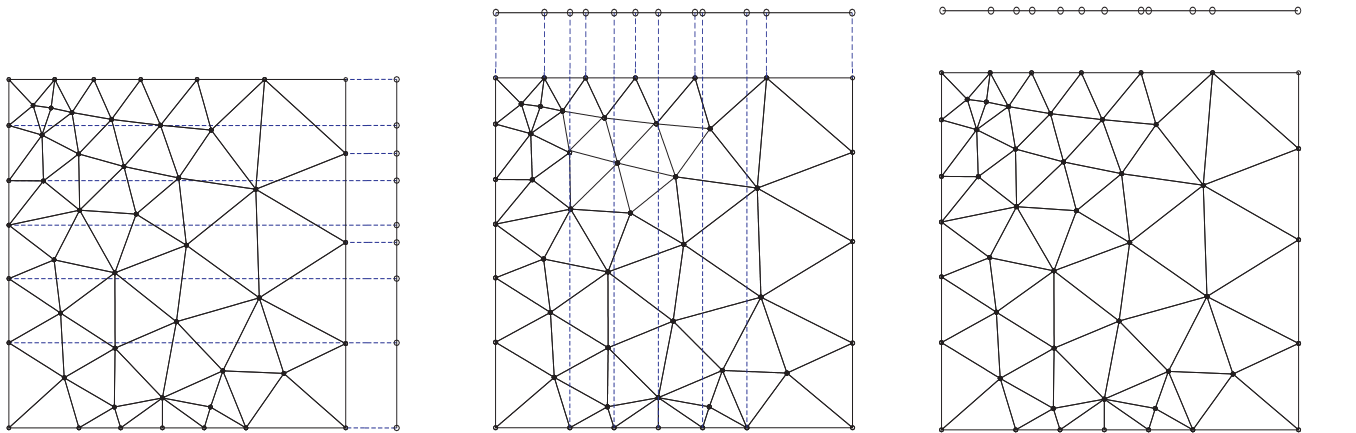
where

$$d_\square(\mathbf{t}, \mathbf{u}) \stackrel{\text{def}}{=} \frac{1}{|\Omega_\square|} \int_{\Gamma_\square^\pm} \mathbf{t} \cdot \llbracket \mathbf{u} \rrbracket \, d\Gamma$$

- **Remark:** Variational framework guarantees Hill-Mandel condition for continuous or discrete \mathbb{U}_\square and \mathbb{T}_\square . \square
- ATS-tensor LARSSON ET AL CMAME 2011

WPBC – Mixed FE-approximation

- Adopted FE: P.w. linear u and p.w. linear (continuous) t in 2D
- Construction of traction mesh

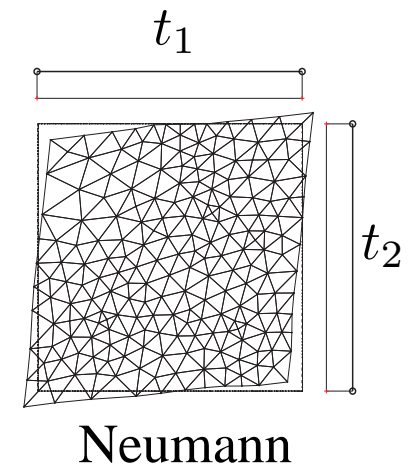
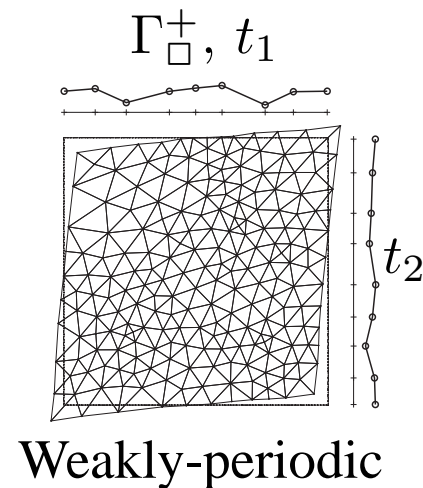
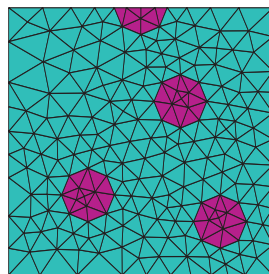


- No need for periodic mesh displacement → Advantageous for adaptive meshing!

WPBC – Mixed FE, cont'd

- **Remarks:**

- Stable combination (in the sense of LBB-condition) for "complete" t -mesh, follows from analysis of contact problems, cf. EL-ABBASI & BATHE 2001, WRIGGERS 2002.
- Stable for arbitrarily coarsened t -mesh; however, $\dim(\mathbb{T}_{\square,h})$ can not be too large
- Special choice: Constant t on each side \leadsto Standard TBC (Neumann b.c)
- Possible to use p.w. constant t from regularity viewpoint



WPBC – Mixed FE, cont'd

- Weak format of periodicity

$$d_{\square}(\delta \mathbf{t}_h, \mathbf{u}_h^s) = 0 \quad \forall \delta \mathbf{t}_h \in \mathbb{T}_{\square, h}$$

- Special case: **Strong periodicity**, $[[\mathbf{u}_h^s]] = \mathbf{0}, \forall \mathbf{X} \in \Gamma_{\square}^+$
 - Increase $\dim(\mathbb{T}_{\square, h})$ *indefinitely*. Not feasible in practice due to instability unless $\mathbb{U}_{\square, h}$ is enlarged indefinitely
 - Introduce *strictly periodic mesh*, i. e. $\text{trace}(\mathbb{U}_{\square, h})$ on Γ_{\square}^+ and $\text{trace}(\mathbb{U}_{\square, h})$ on Γ_{\square}^- are identical. Choose

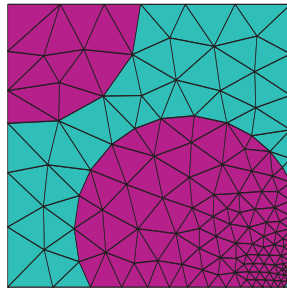
$$\mathbb{T}_{\square, h} = \text{trace}(\mathbb{U}_{\square, h})_{\Gamma_{\square}^+} = \text{trace}(\mathbb{U}_{\square, h})_{\Gamma_{\square}^-}$$

Setting $\delta \mathbf{t}_h = [[\mathbf{u}_h^s]] \in \mathbb{T}_{\square, h} \Rightarrow [[\mathbf{u}_h^s]] = \mathbf{0}$, i. e. strong periodicity!

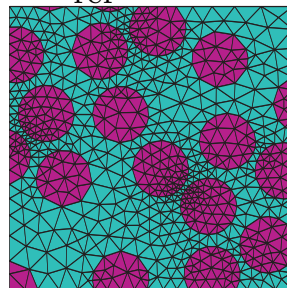
- **Note:** Strong periodicity does not require any particular arrangement or precautions. It is merely considered as a special choice out of many possible meshes.

WPBC – Typical results

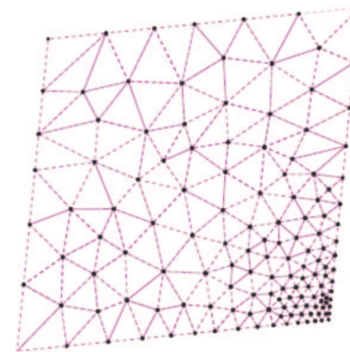
- Snapshots of deformed RVEs for single realization, from CMAME, in print



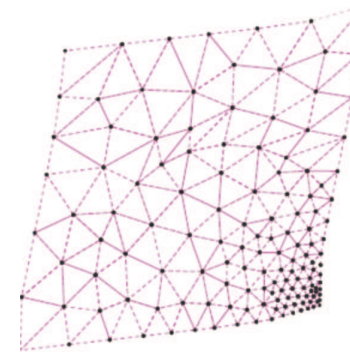
$$\frac{L_{\square}}{L_{\text{ref}}} = 1.25$$



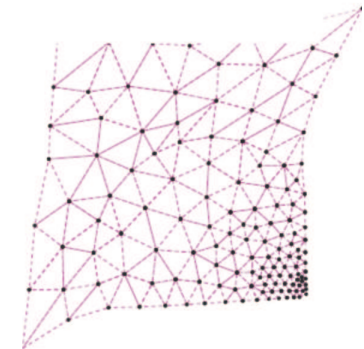
$$\frac{L_{\square}}{L_{\text{ref}}} = 5.00$$



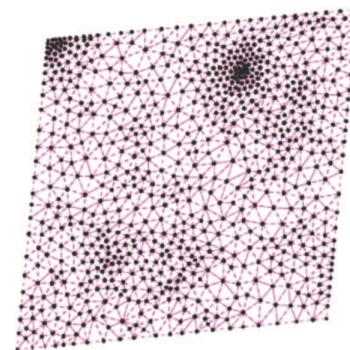
Dirichlet



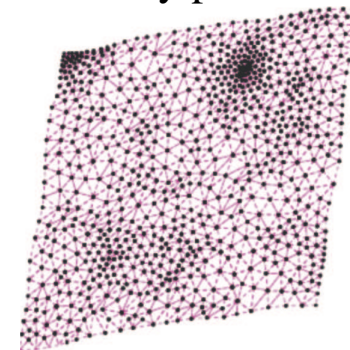
Weakly periodic



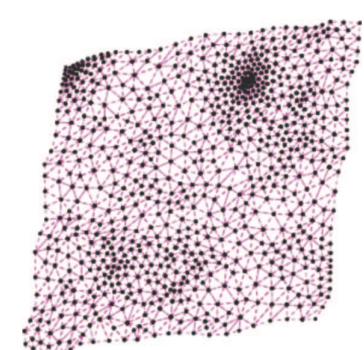
Neumann



Dirichlet



Weakly periodic



Neumann

- **Remark:** Non-periodic mesh for weakly periodic bc! Suitable for adaptive meshing.

SVE-problems

- Displacement b.c. (DBC), Traction b.c. (TBC), Weak periodicity b.c. (WPBC)
- Macroscale data (input or output) for SVE-problem:
 - Macrostrain $\bar{H} \in \mathbb{R}^{3 \times 3}$ or Macrostress $\bar{P} \in \mathbb{R}^{3 \times 3}$
- Subscale variables (fields)
 - Displacement $u \in U_{\square}^D, U_{\square}^P, U_{\square}^N$
 - Tractions $t \in T_{\square}$
- Classical versions: DBC with strain control, TBC with stress control

	DBC	WPBC	TBC
Strain control \bar{H}	u \bar{P} from post-proc	u, t \bar{P} from post-proc	u, \bar{P}
Stress control \bar{P}	u, \bar{H}	u, t, \bar{H}	u \bar{H} from post-proc