

Three-Phase Displacement Theory: Hyperbolic Models and Analytical Solutions

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Contents

Discussion of general conditions for relative permeabilities to ensure hyperbolicity

Presentation of the analytical solution to the Riemann problem

Implementation in a front-tracking method

Streamline simulation results

Displacement Theory

Assumptions:

- One-dimensional flow
- Immiscible fluids
- Incompressible fluids
- Homogeneous rigid porous medium
- Multiphase flow extension of Darcy's law
- Gravity and capillarity are not considered
- Constant fluid viscosities

Displacement Theory cont'd

Mass conservation for each phase:

$$\partial_t(m_\alpha) + \partial_x(F_\alpha) = 0, \quad \alpha = w, g, o$$

$$m_\alpha = \rho_\alpha S_\alpha \phi$$

$$F_\alpha = -\rho_\alpha k \lambda_\alpha \partial_x p$$

Saturations add up to one:

$$\sum_{\alpha=w,g,o} S_\alpha \equiv 1$$

Displacement Theory cont'd

If the fractional flow approach is used:

- “Pressure equation”

$$\partial_x(v_T) = 0$$

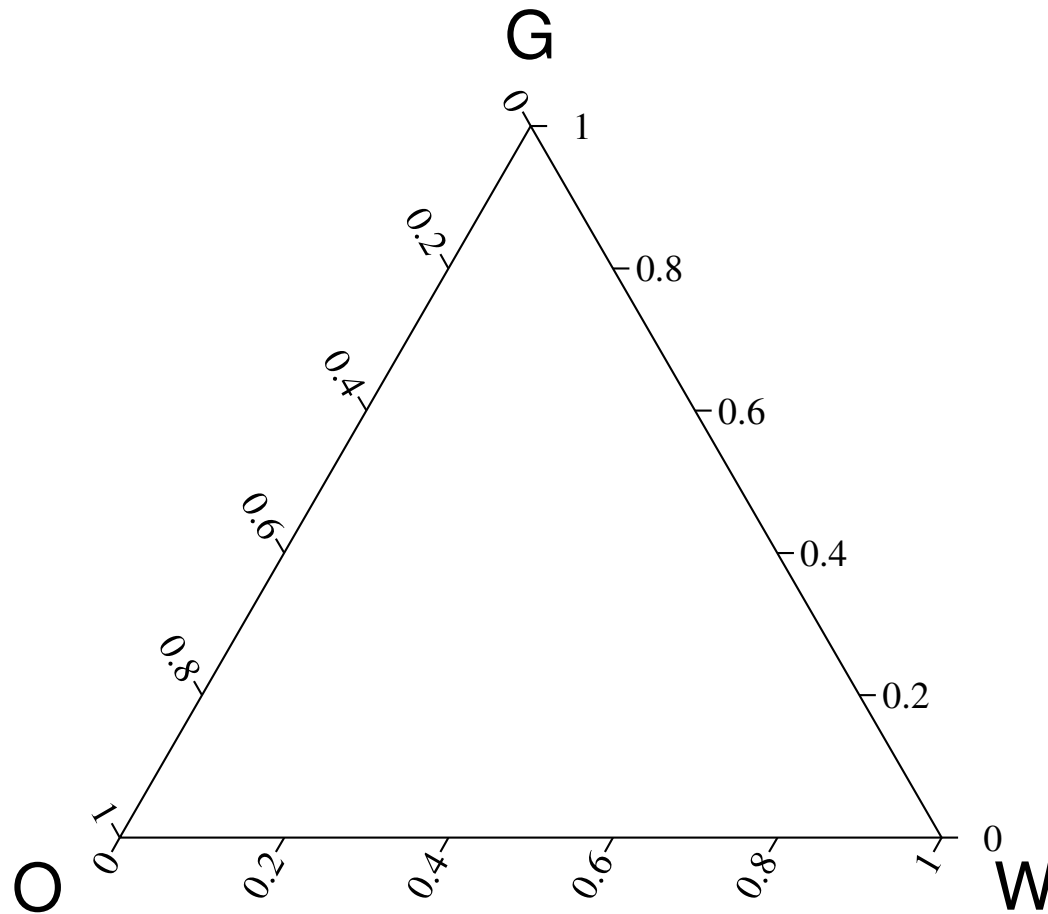
$$v_T = -\frac{1}{\phi}k\lambda_T\partial_x p$$

- System of “saturation equations”

$$\partial_t \begin{pmatrix} S_w \\ S_g \end{pmatrix} + v_T \partial_x \begin{pmatrix} f_w \\ f_g \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

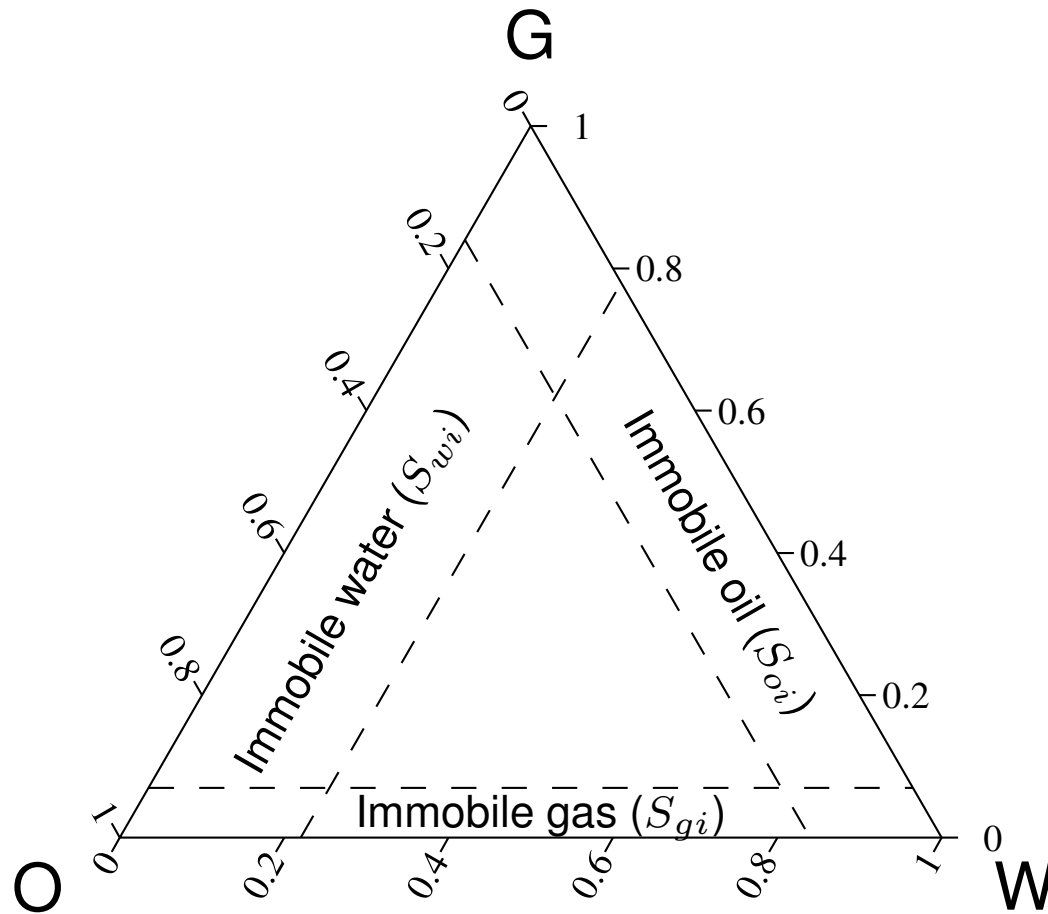
Displacement Theory cont'd

Saturation triangle:



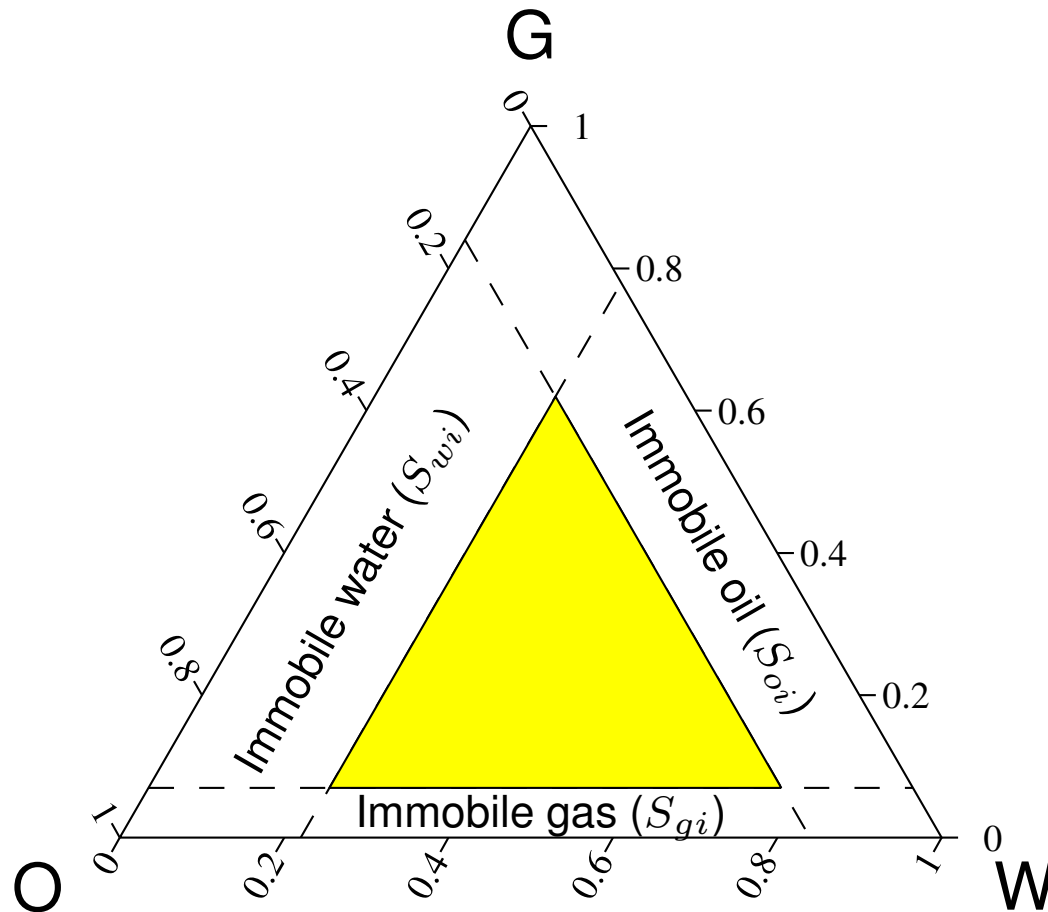
Displacement Theory cont'd

Saturation triangle:



Displacement Theory cont'd

Saturation triangle:



Reduced saturations:

$$\tilde{S}_\alpha := \frac{S_\alpha - S_{\alpha i}}{1 - \sum_{\beta=1}^3 S_{\beta i}}$$



Renormalized triangle

Character of the System

The character of the system

$$\partial_t \begin{pmatrix} u \\ v \end{pmatrix} + v_T \partial_x \begin{pmatrix} f \\ g \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \Leftrightarrow \partial_t \mathbf{u} + v_T \partial_x \mathbf{f} = \mathbf{0}$$

is determined by the eigenvalues (ν_1, ν_2) and eigenvectors $(\mathbf{r}_1, \mathbf{r}_2)$ of the *Jacobian matrix*:

$$\mathbf{A}(\mathbf{u}) := \mathbf{D}_u \mathbf{f} = \begin{pmatrix} f_{,u} & f_{,v} \\ g_{,u} & g_{,v} \end{pmatrix}$$

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Hyperbolic: The eigenvalues are *real* and the Jacobi matrix is *diagonalizable*. Strictly hyperbolic: distinct eigenvalues $\nu_1 < \nu_2$.

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Hyperbolic: The eigenvalues are *real* and the Jacobi matrix is *diagonalizable*. Strictly hyperbolic: distinct eigenvalues $\nu_1 < \nu_2$.

Elliptic: The eigenvalues are *complex*.

Character of the System cont'd

It was long believed that the system was strictly hyperbolic for *any* set of relative permeability functions

Character of the System cont'd

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However: existence of elliptic regions proved 1985–95

- Bell et al., Fayers, Guzmán and Fayers, Hicks and Grader, ..
- Shearer and Trangenstein, Holden et al, ...

Character of the System cont'd

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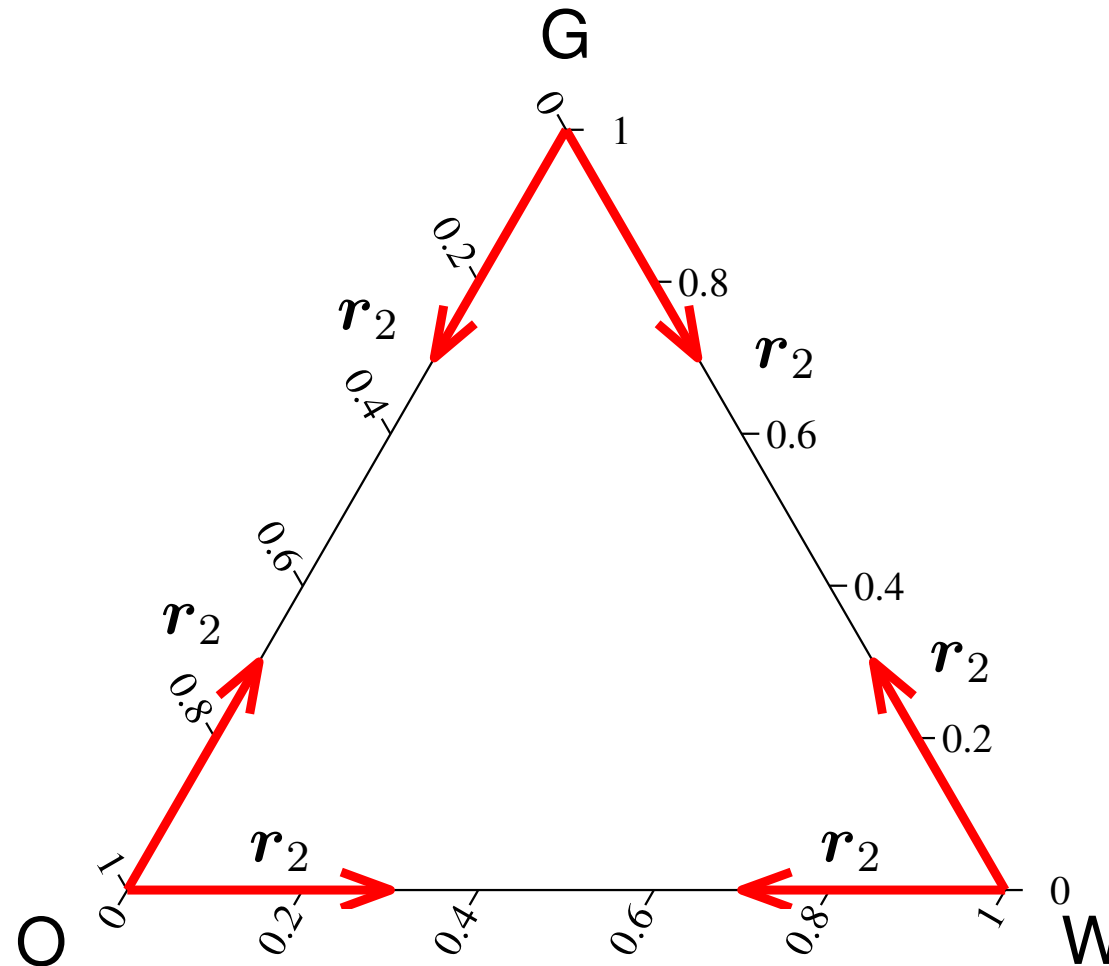
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Approach in the existing literature:

- **Assume** “reasonable” conditions for relative permabilities on the edges
 - “Zero-derivative” conditions
 - “Interaction” conditions
- **Infer** loss of strict hyperbolicity inside the saturation triangle

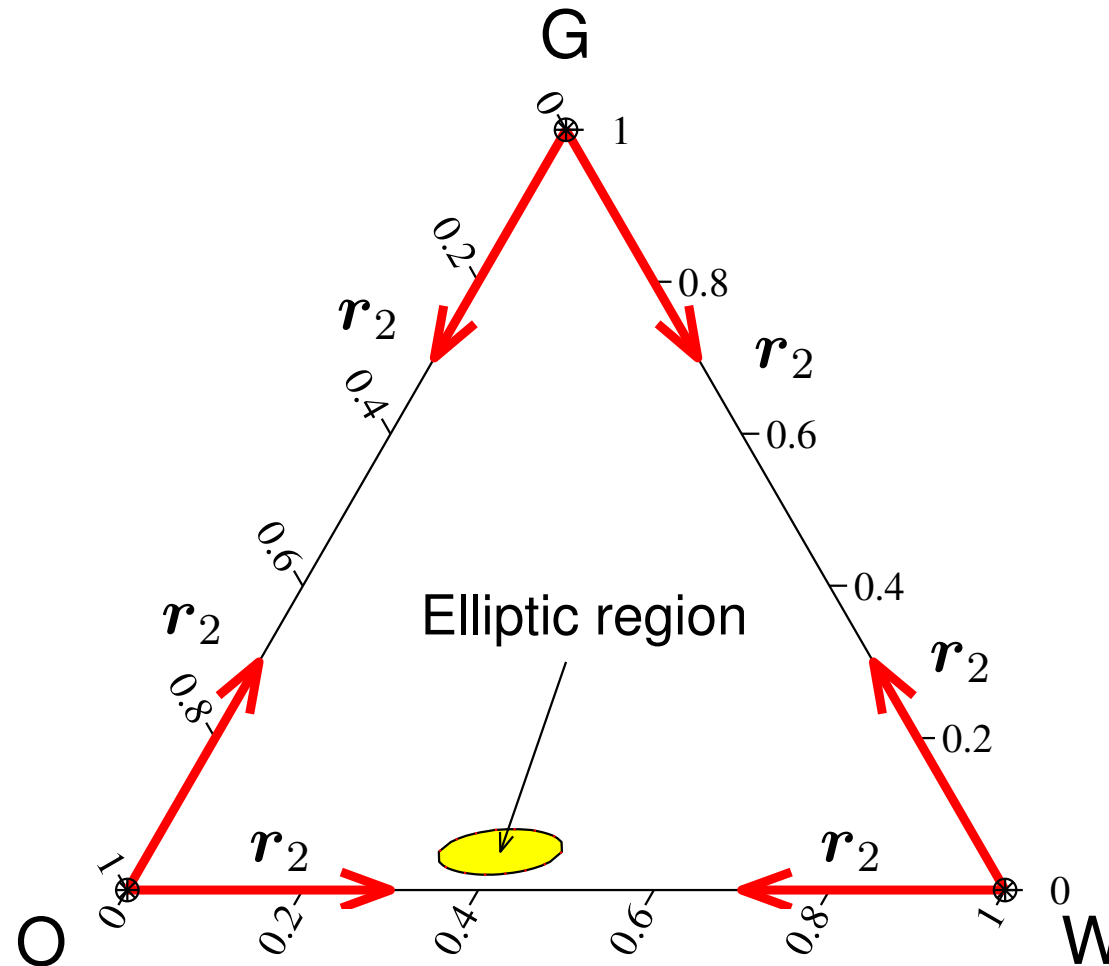
Character of the System cont'd

Traditional assumed behavior of fast eigenvectors (r_2)



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Traditional assumed behavior of fast eigenvectors (r_2)



Character of the System cont'd

Consequences of ellipticity of the system:

- Flow depends on *future* boundary conditions
- The solution is *unstable*: arbitrarily close initial and injected saturations yield nonphysical oscillatory waves

Character of the System cont'd

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However:

- The elliptic region can be shrunk to an **umbilic point** only if interaction between phases is ignored:

$$k_{r\alpha} = k_{r\alpha}(S_\alpha), \quad \alpha = 1, \dots, 3$$

- This model is *not* supported by experiments and pore-scale physics
- Umbilic points still act as “repellers” for classical waves

Relative Permeabilities

Juanes and Patzek – New approach:

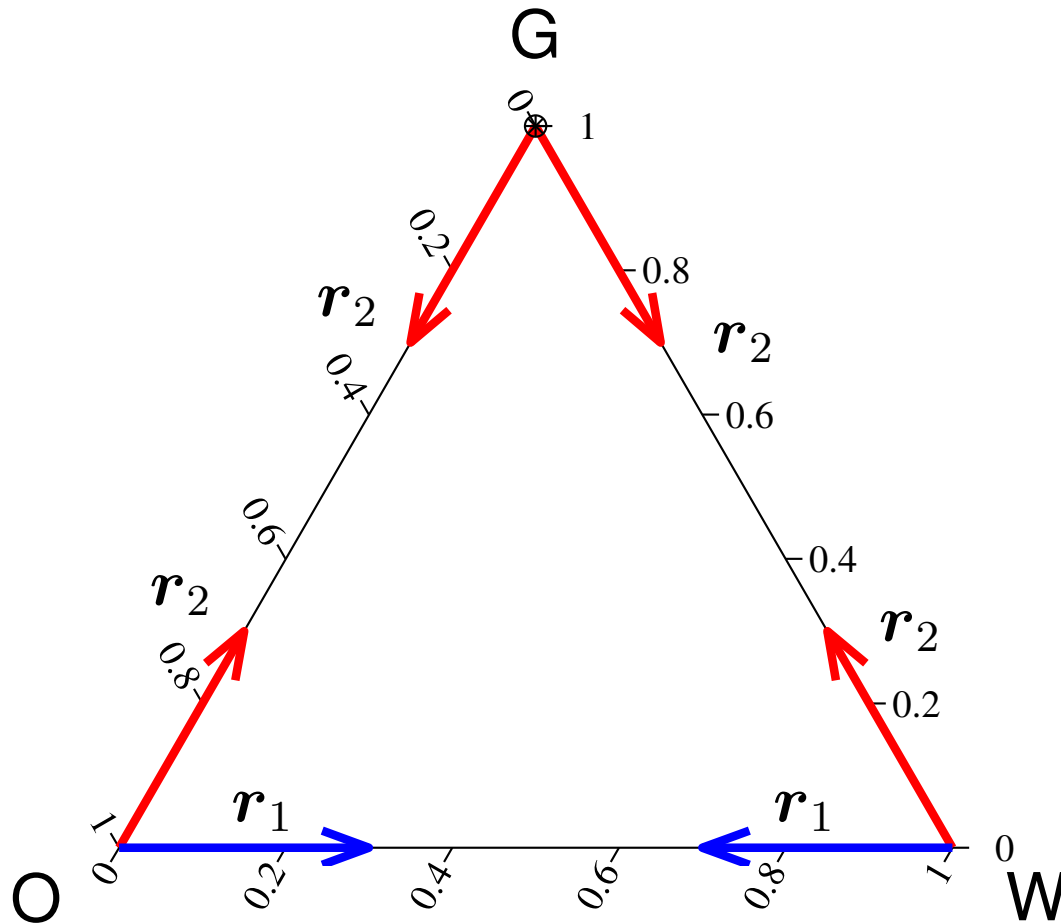
- **Assume** the system is strictly hyperbolic
- **Infer** conditions on relative permeabilities

Key observation:

- Whenever gas is present as a continuous phase, its mobility is much higher than that of the other two fluids
- Fast paths \longleftrightarrow changes in gas saturation

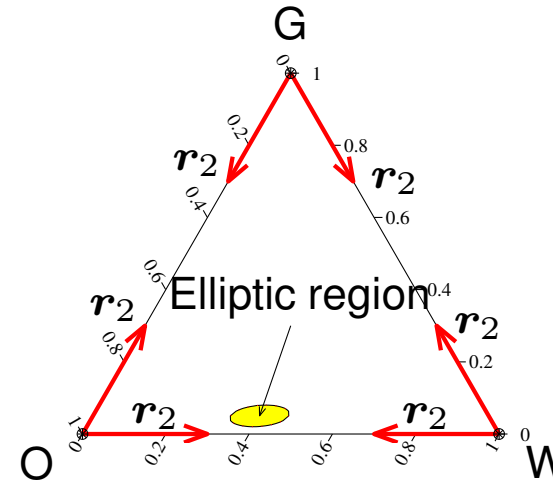
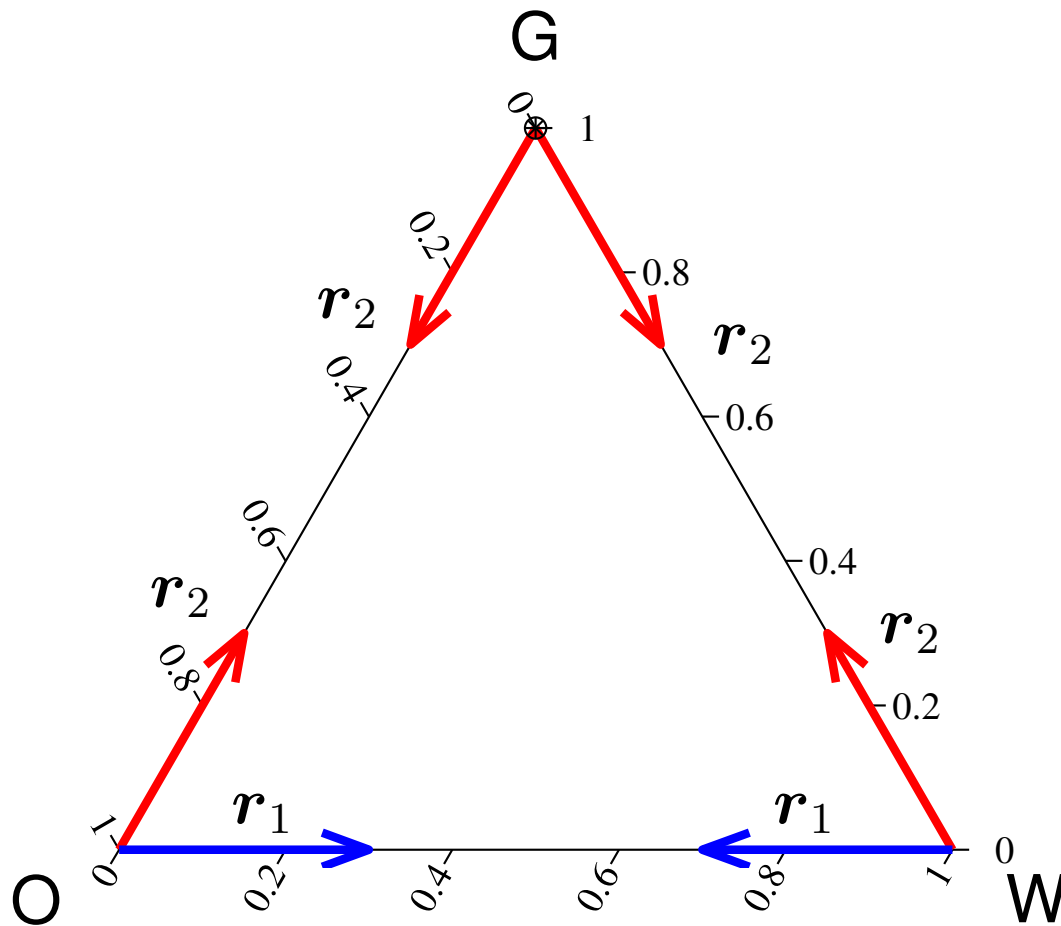
Relative Permeabilities

Proposed behavior of eigenvectors (r_1 , r_2)



Relative Permeabilities

Proposed behavior of eigenvectors (r_1 , r_2)



Relative Permeabilities cont'd

Two types of conditions:

- **Condition I.** Eigenvectors are parallel to each edge
- **Condition II.** Strict hyperbolicity along each edge

In particular, on the **OW edge**:

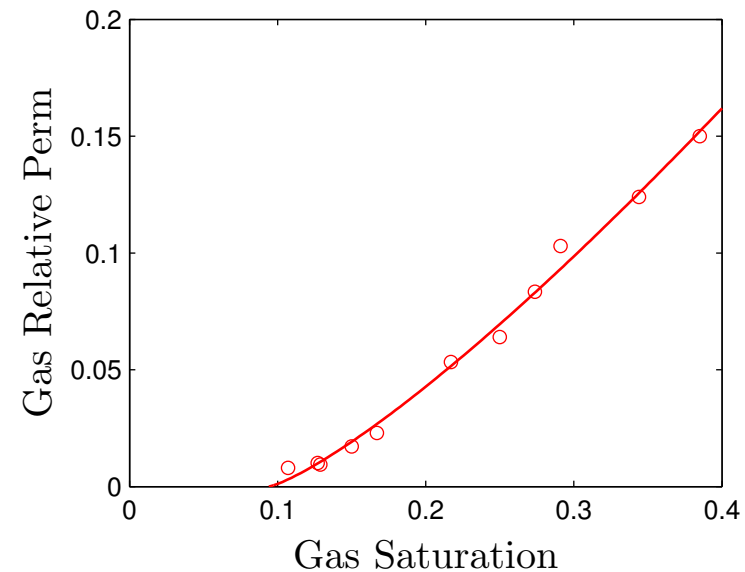
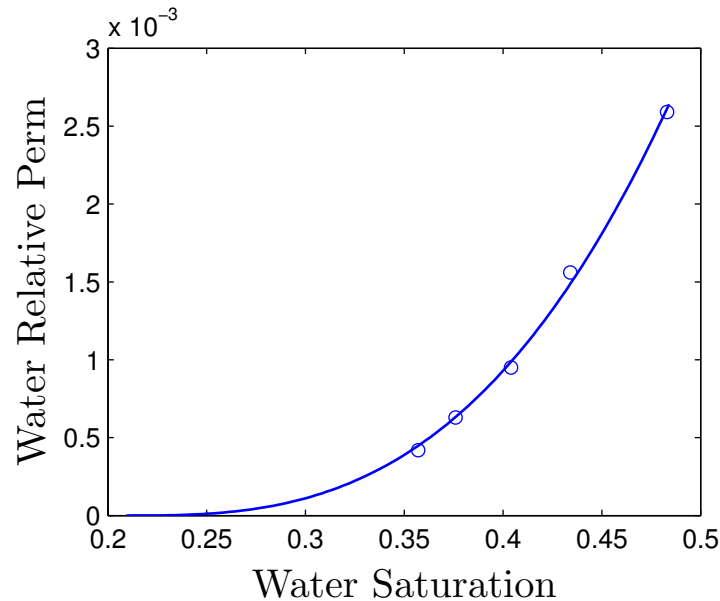
Condition	Frac. flows	Mobilities
I	$g_{,u} = 0$	$\Leftrightarrow \lambda_{g,u} = 0$
II	$g_{,v} - f_{,u} > 0$	$\Leftrightarrow \lambda_{g,v} > \lambda_{w,u} - \lambda_{T,u} \frac{\lambda_w}{\lambda_T}$

Condition II requires that the gas relative permeability has a *positive derivative* at its endpoint saturation.

Relative Permeabilities cont'd

Remarks:

- Necessary condition for strict hyperbolicity
- Can be justified from pore-scale physics (bulk flow vs. corner flow)
- Supported by experimental data (Oak's steady-state)



Relative Permeabilities cont'd

A simple model:

$$k_{rw}(u) = u^2$$

$$k_{rg}(v) = (\beta_g v + (1 - \beta_g)v^2), \quad \beta_g > 0$$

$$k_{ro}(u, v) = (1 - u - v)(1 - u)(1 - v)$$

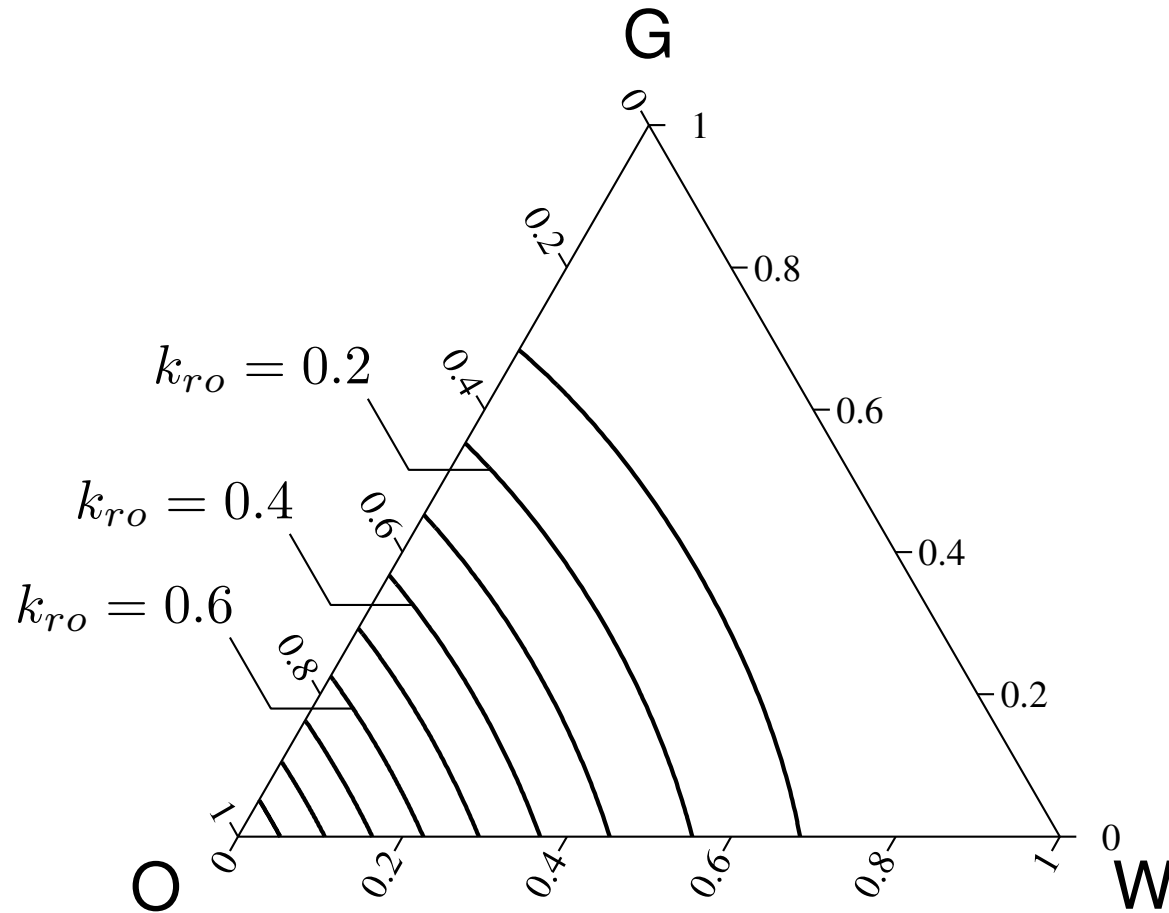
with reasonable values of viscosities:

$$\mu_w = 1, \quad \mu_g = 0.03, \quad \mu_o = 2 \text{ cp}$$

and a small value of the endpoint slope: $\beta_g = 0.1$

Relative Permeabilities cont'd

Oil isoperms:



Analytical Solution

Riemann problem: Find a weak solution to the 2×2 system

$$\partial_t \mathbf{u} + v_T \partial_x \mathbf{f} = \mathbf{0}, \quad -\infty < x < \infty, \quad t > 0$$

$$\mathbf{u}(x, 0) = \begin{cases} \mathbf{u}_l & \text{if } x < 0 \\ \mathbf{u}_r & \text{if } x > 0 \end{cases}$$

Previous work:

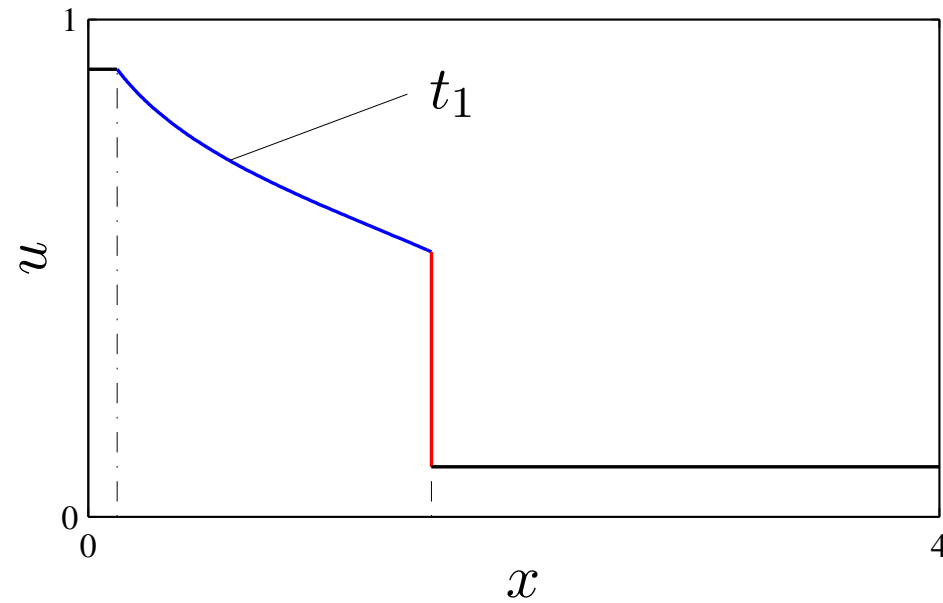
- Sequence of two successive two-phase displacements (Kyte et al., Pope, ..)
- Triangular systems (Gimse et al., ..)

New results by Juanes and Patzek:

- *A complete classification all wave types*
- Solution of Riemann problem (structure of waves)

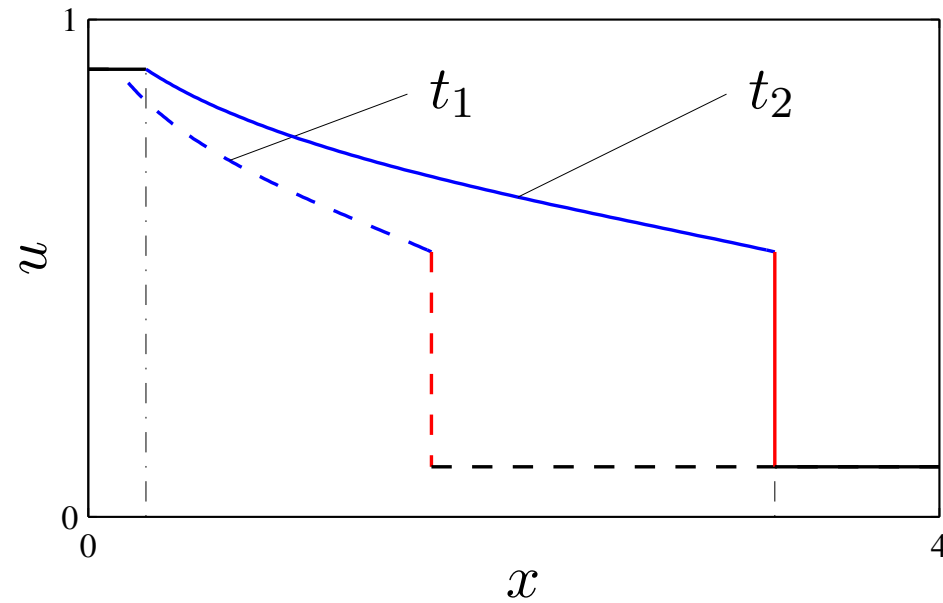
Analytical Solution cont'd

Self-similarity (“stretching”, “coherence”):



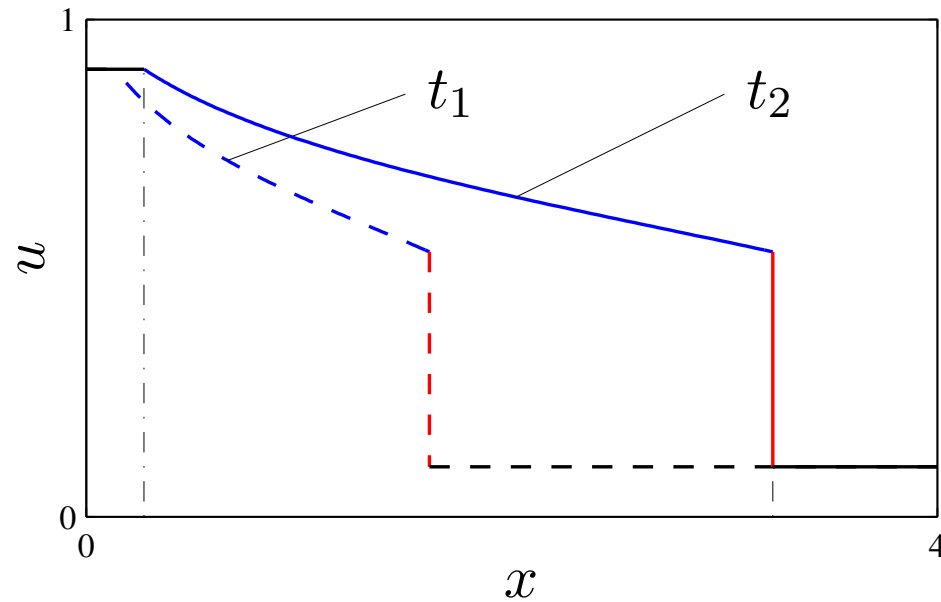
Analytical Solution cont'd

Self-similarity (“stretching”, “coherence”):



Analytical Solution cont'd

Self-similarity (“stretching”, “coherence”):



$$u(x, t) = U(\zeta), \quad \zeta := \frac{x}{\int_0^t v_T(\tau) d\tau}$$

Analytical Solution cont'd

Using self-similarity, the Riemann problem is a boundary value problem:

$$(A(U) - \zeta I)U' = \mathbf{0}, \quad -\infty < \zeta < \infty$$

with boundary conditions

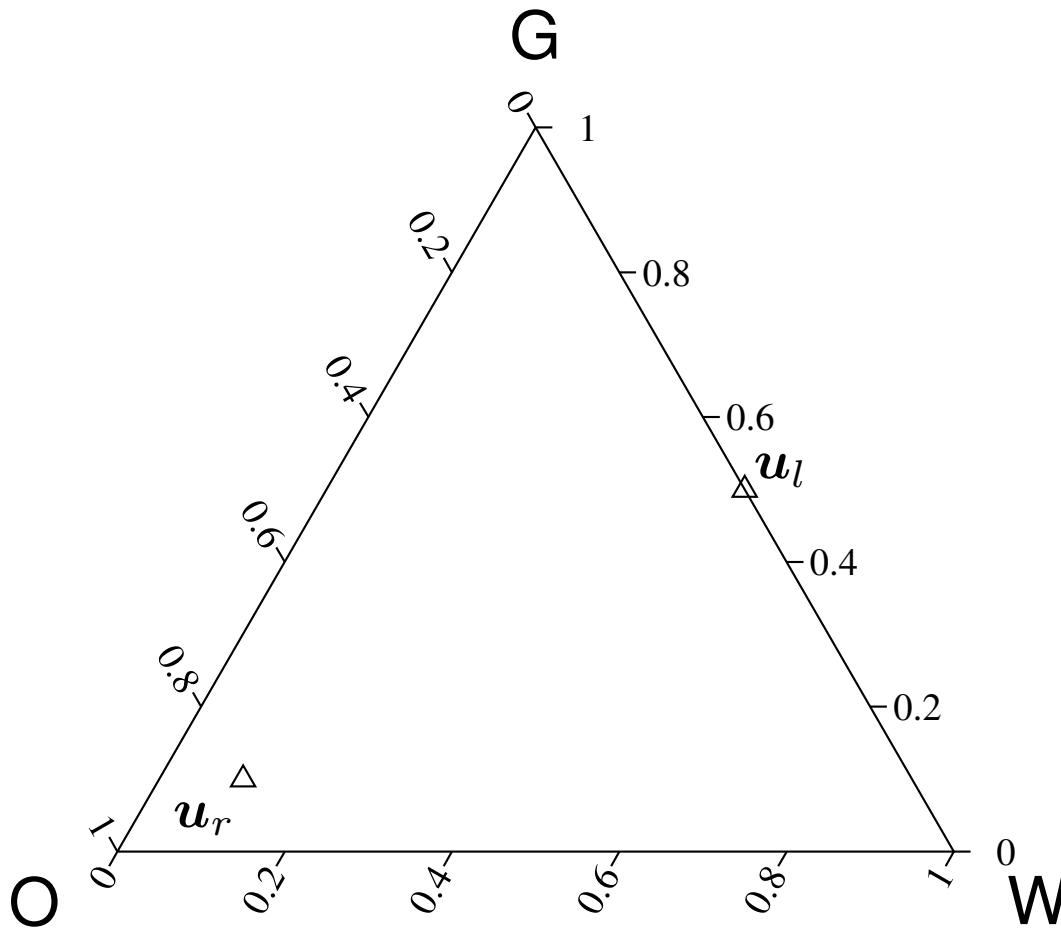
$$U(-\infty) = \mathbf{u}_l, \quad U(\infty) = \mathbf{u}_r$$

Strict hyperbolicity \longrightarrow wave separation:

$$\mathbf{u}_l \xrightarrow{\mathcal{W}_1} \mathbf{u}_m \xrightarrow{\mathcal{W}_2} \mathbf{u}_r$$

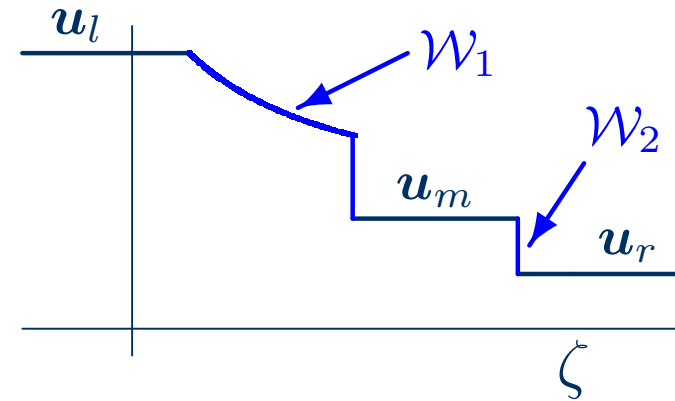
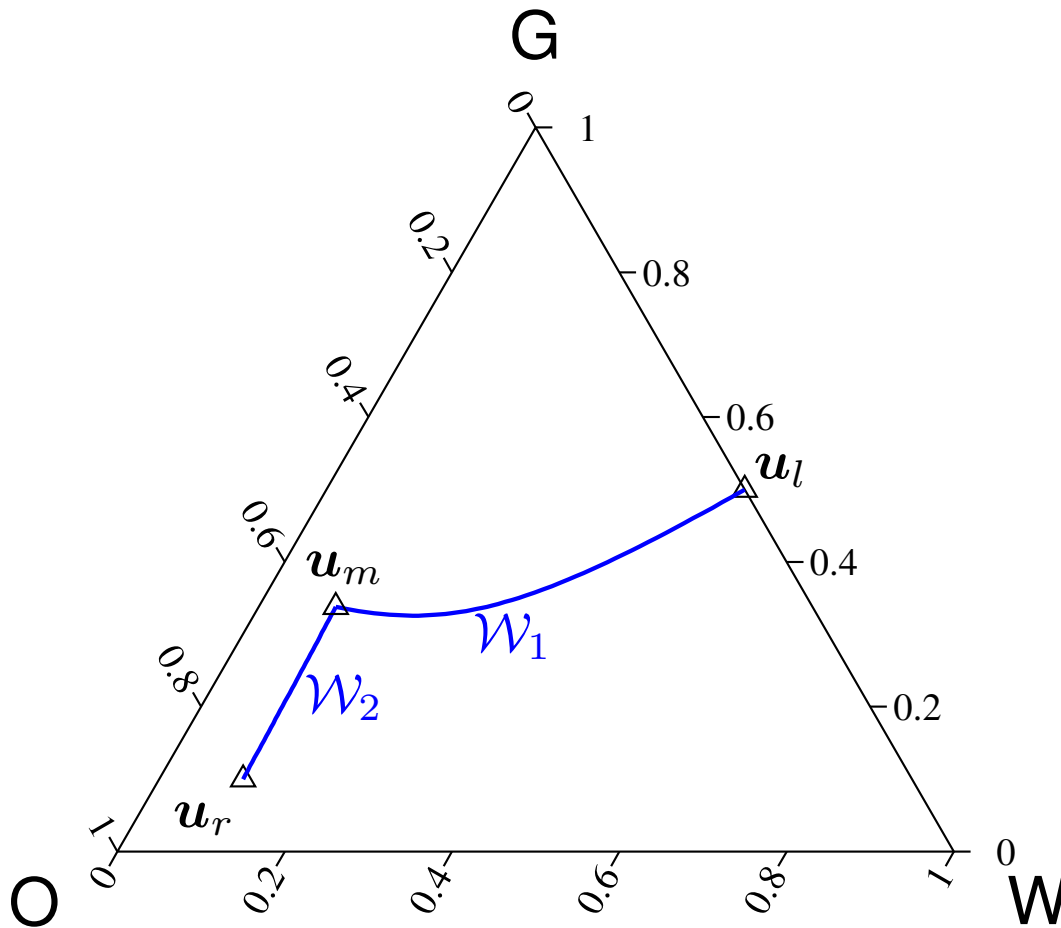
Analytical Solution cont'd

Schematic of Riemann solution



Analytical Solution cont'd

Schematic of Riemann solution



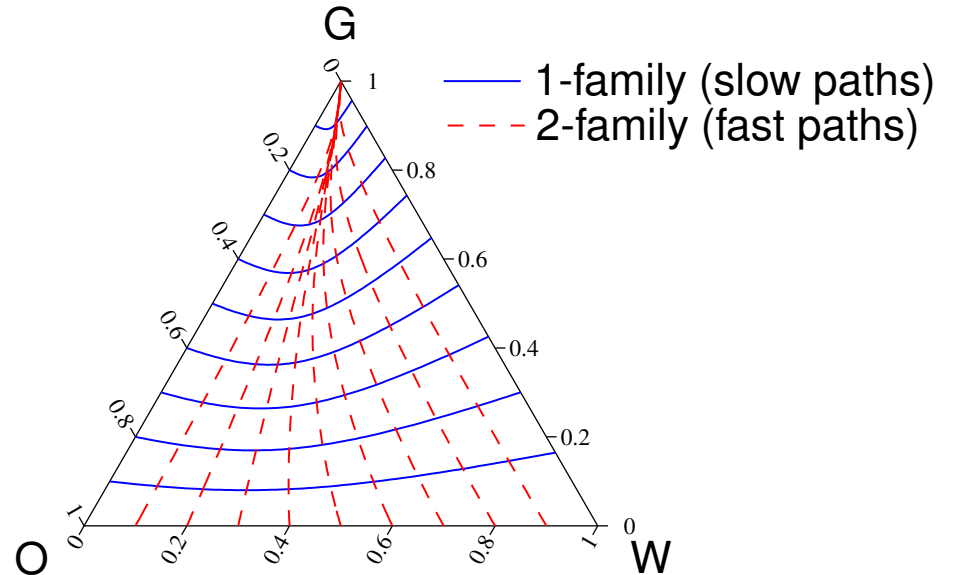
Wave Structure: Rarefactions

If the solution is smooth, it satisfies:

$$(A(U) - \zeta I)U' = 0$$

\nearrow \nwarrow
 eigenvalue (ν_p) eigenvector (r_p)

A smooth solution (**rarefaction**) must lie on a curve whose tangent is in the direction of the eigenvector (**integral curve**)



Wave Structure: Rarefactions cont'd

Admissibility of a rarefaction wave

- To avoid a multiple-valued solution, ν_p must **increase monotonically** along the curve

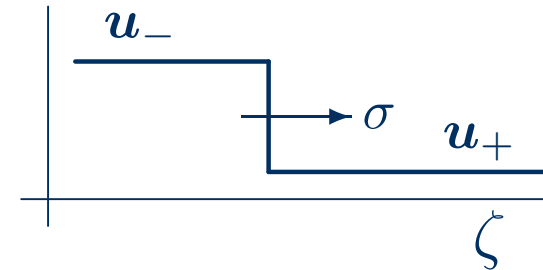
$$\mathbf{u}_l \xrightarrow{\mathcal{R}_p} \mathbf{u}_r$$

- Thus, rarefaction curves \mathcal{R}_p are **subsets** of integral curves \mathcal{I}_p

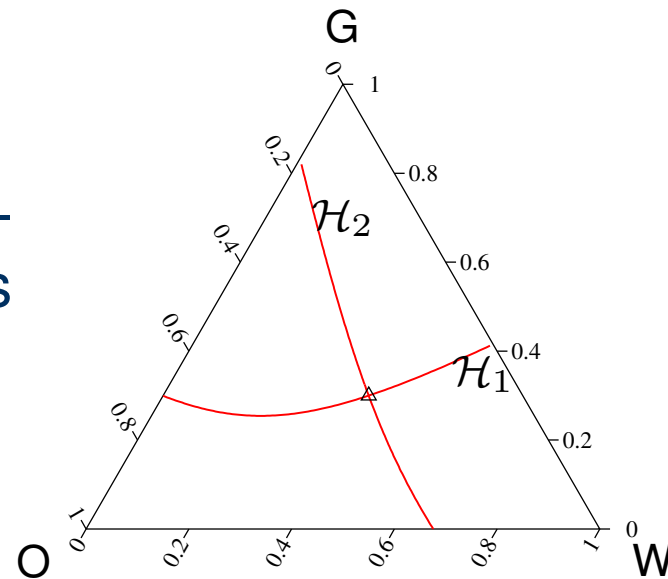
Wave Structure: Shocks

If the solution is discontinuous, it must satisfy the **Rankine-Hugoniot jump condition**:

$$f(u_+) - f(u_-) = \sigma \cdot [u_+ - u_-]$$



The set of states which can be connected satisfying the jump condition is called the **Hugoniot locus**



Wave Structure: Shocks cont'd

Admissibility of a shock wave

- Not every discontinuity satisfying the jump condition is a valid shock
- Characteristics of the p -family must go **into** the shock (**Lax entropy condition**):

$$\nu_p(\mathbf{u}_-) > \sigma_p > \nu_p(\mathbf{u}_+),$$

- Thus, shock curves \mathcal{S}_p are **subsets** of Hugoniot loci \mathcal{H}_p

Wave Structure: Rarefaction-Shocks

Genuine nonlinearity: eigenvalues vary **monotonically** along integral curves

Wave Structure: Rarefaction-Shocks

Genuine nonlinearity: eigenvalues vary **monotonically** along integral curves

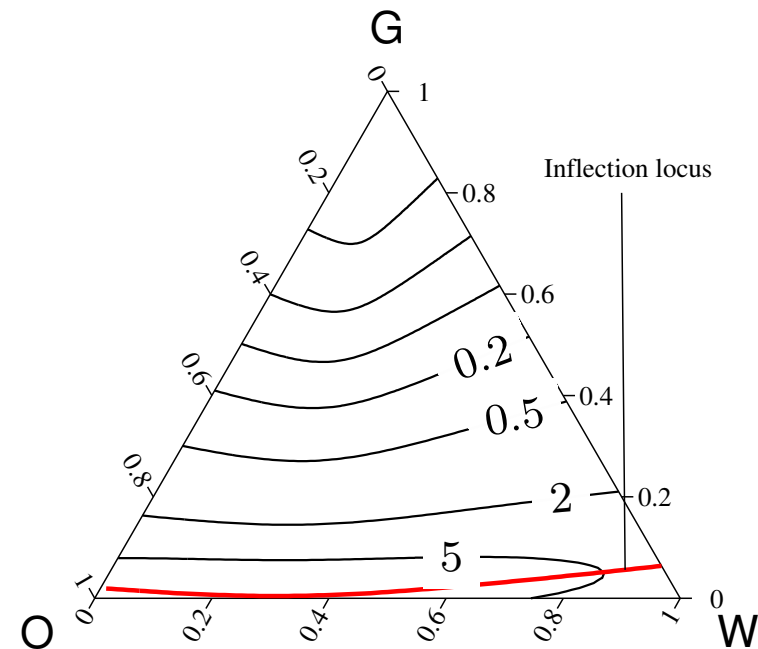
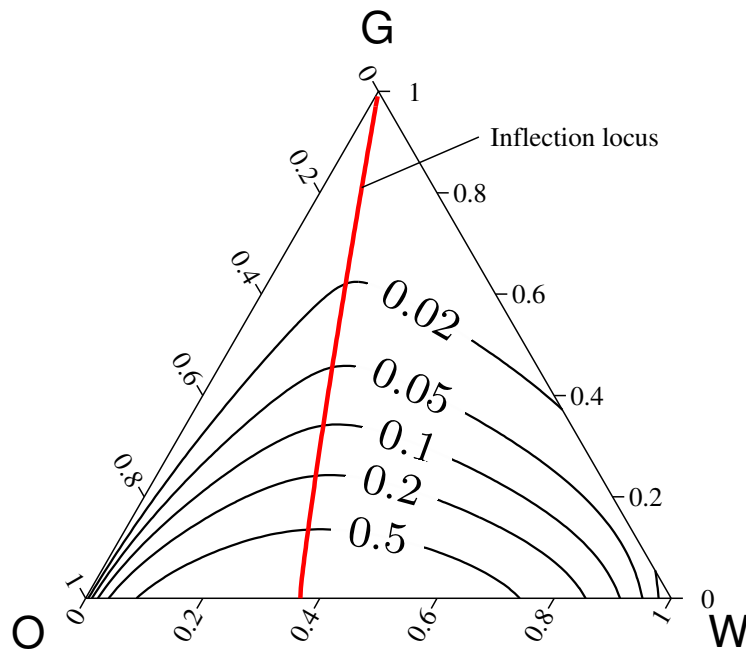
This is **not** the case in multiphase flow, where each wave may involve rarefactions **and** shocks

Wave Structure: Rarefaction-Shocks

Genuine nonlinearity: eigenvalues vary **monotonically** along integral curves

This is **not** the case in multiphase flow, where each wave may involve rarefactions **and** shocks

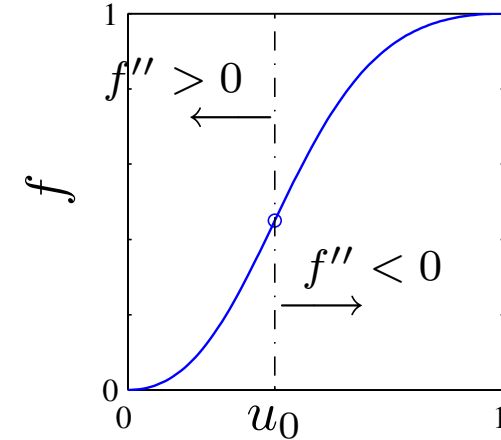
Inflection locus: set of points at which eigenvalues attain a local maximum when moving along integral curves



Wave Structure: Rarefaction-Shocks cont'd

Properties of the inflection loci:

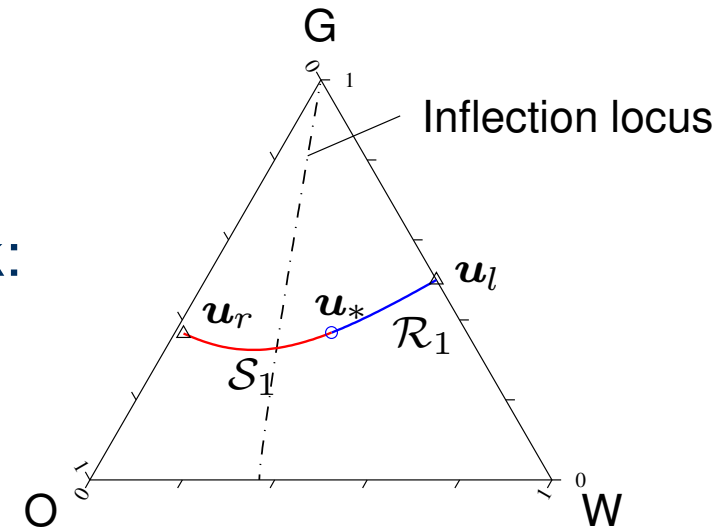
- **Single** curves, transversal to integral curves
- Correspond to **maxima** of eigenvalues



Consequences:

- **At most** one rarefaction and one shock
- Rarefaction is **always slower** than shock:

$$u_l \xrightarrow{\mathcal{R}_p} u_* \xrightarrow{\mathcal{S}_p} u_r$$



Wave Structure: Rarefaction-Shocks cont'd

Admissibility of a rarefaction-shock wave

- Eigenvalue ν_p must increase monotonically along the curve

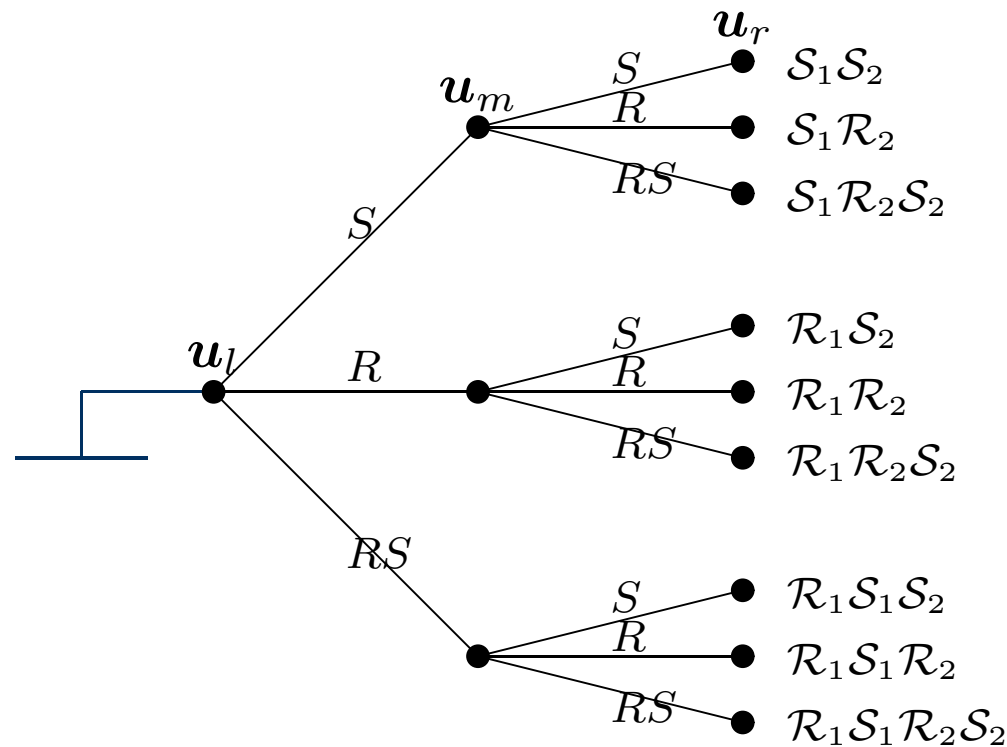
$$\mathbf{u}_l \xrightarrow{\mathcal{R}_p} \mathbf{u}_*$$

- The shock must satisfy the **extended-Lax entropy condition**:

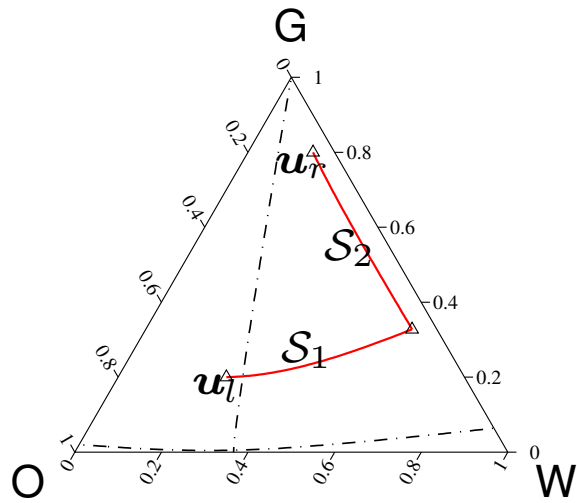
$$\nu_p(\mathbf{u}_*) = \sigma_p > \nu_p(\mathbf{u}_r)$$

Wave Structure cont'd

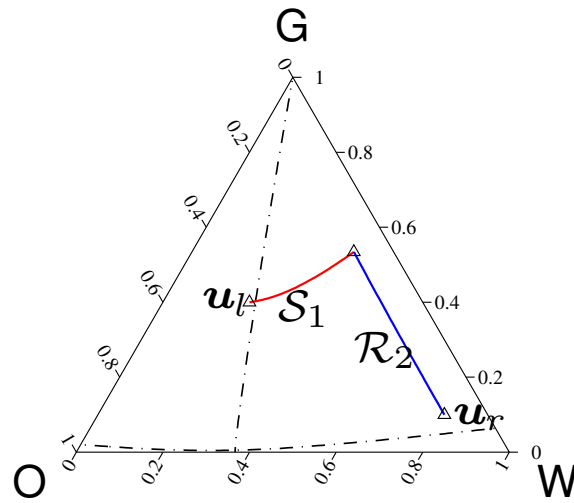
Complete set of solutions: 9 different wave configurations



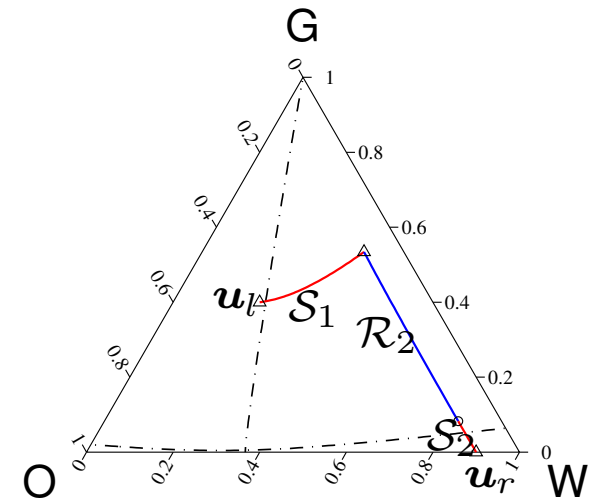
Complete Set of Solutions



(a) $S_1 S_2$

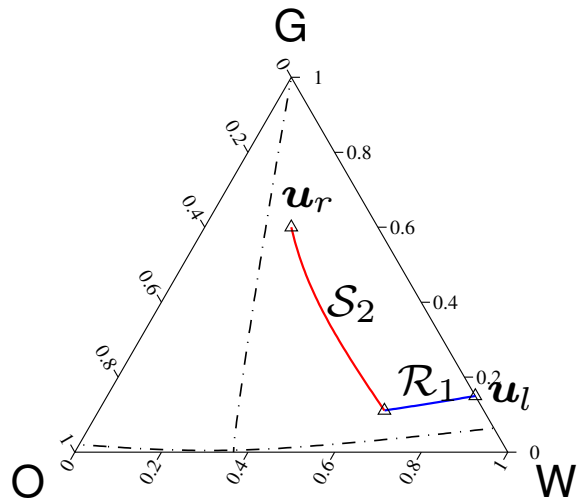


(b) $S_1 R_2$

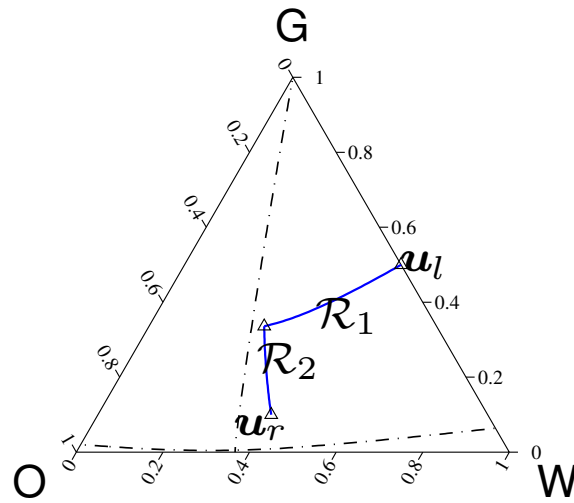


(c) $S_1 R_2 S_2$

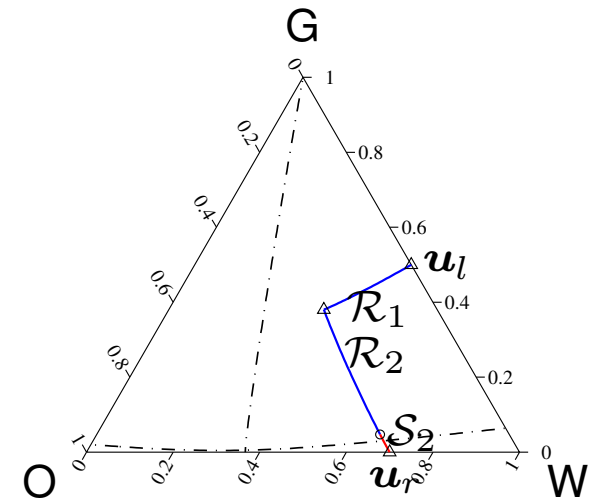
Complete Set of Solutions cont'd



(d) $\mathcal{R}_1 S_2$

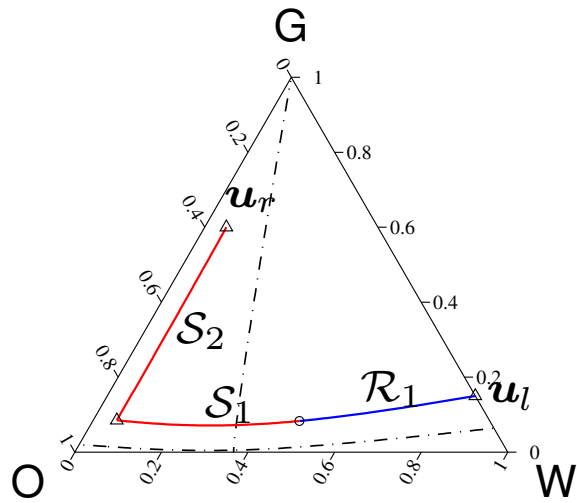


(e) $\mathcal{R}_1 \mathcal{R}_2$

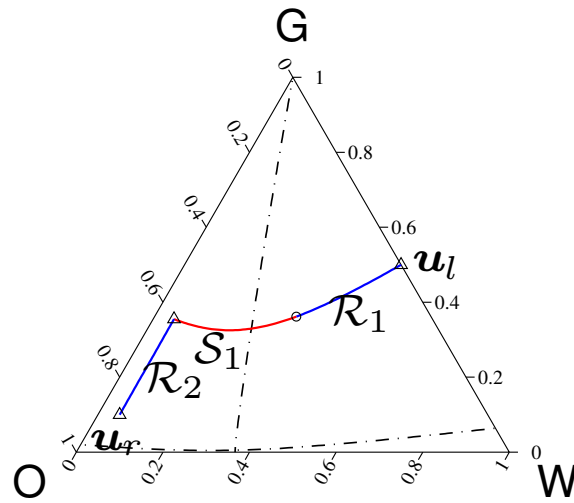


(f) $\mathcal{R}_1 \mathcal{R}_2 S_2$

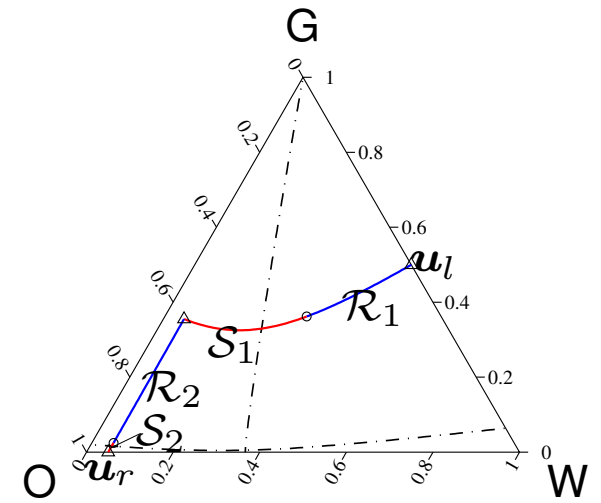
Complete Set of Solutions cont'd



(g) $R_1S_1S_2$



(h) $R_1S_1R_2$



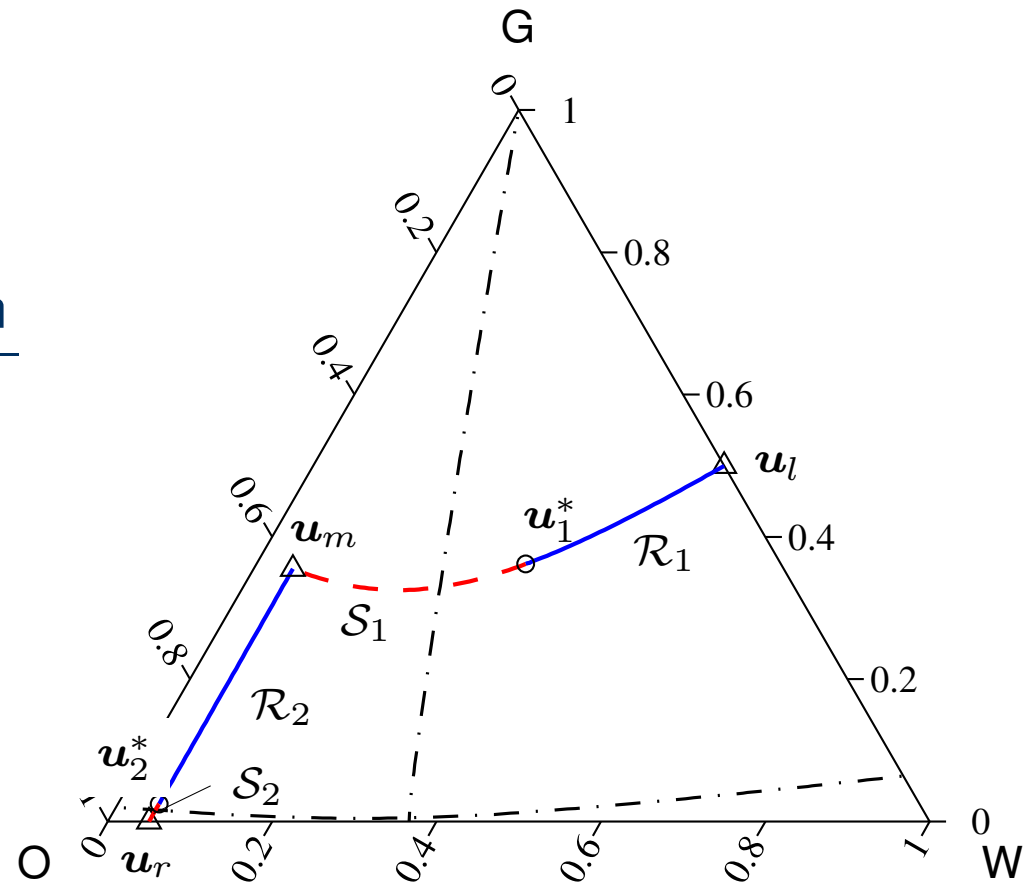
(i) $R_1S_1R_2S_2$

Example 1

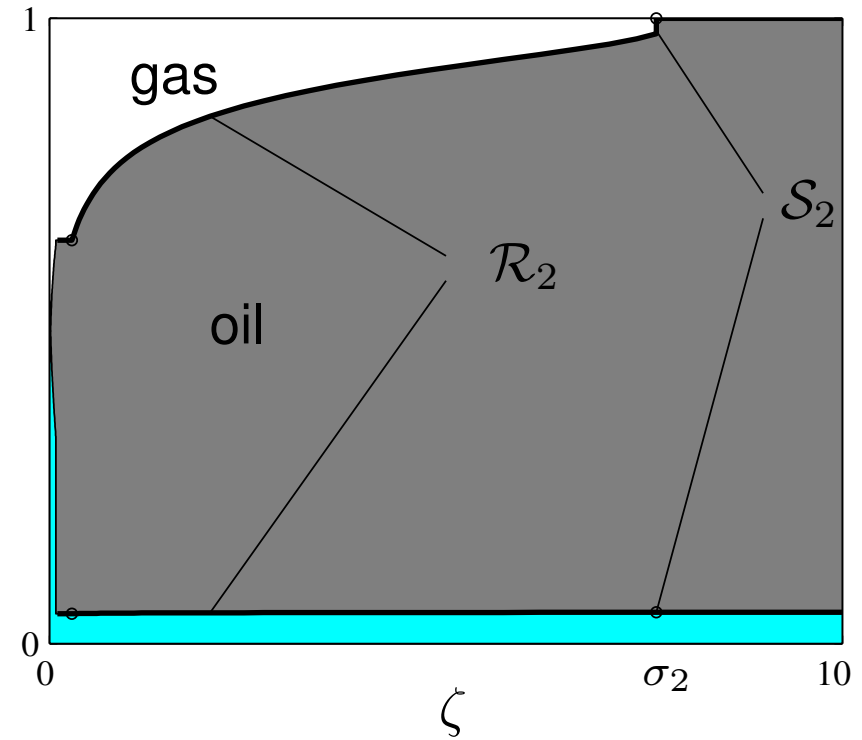
Injection of water and gas into an oil-filled core (with some mobile water)

Problem of great practical interest

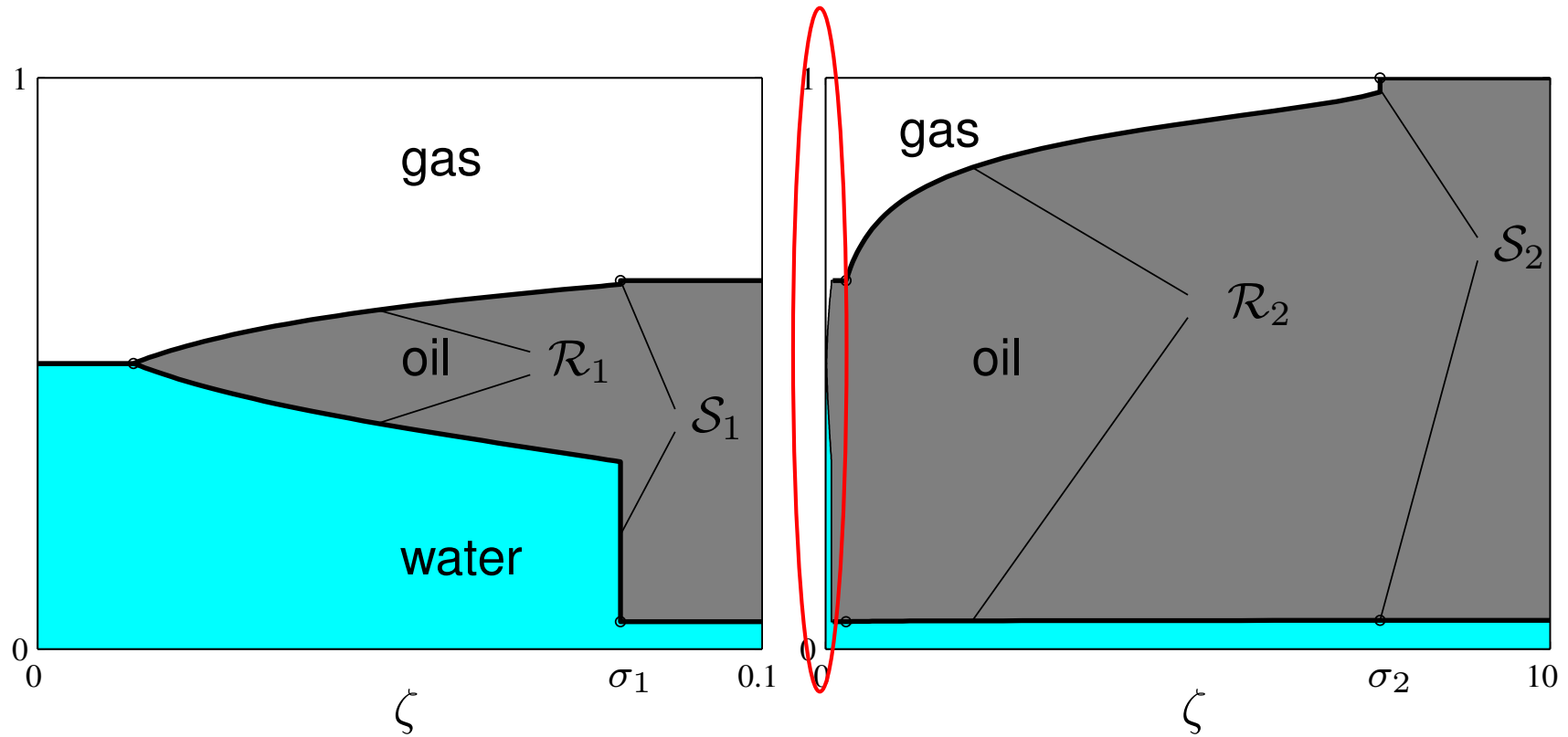
Injected saturation	Initial saturation
$S_g = 0.5$	$S_g = 0$
$S_o = 0$	$S_o = 0.95$
$S_w = 0.5$	$S_w = 0.05$



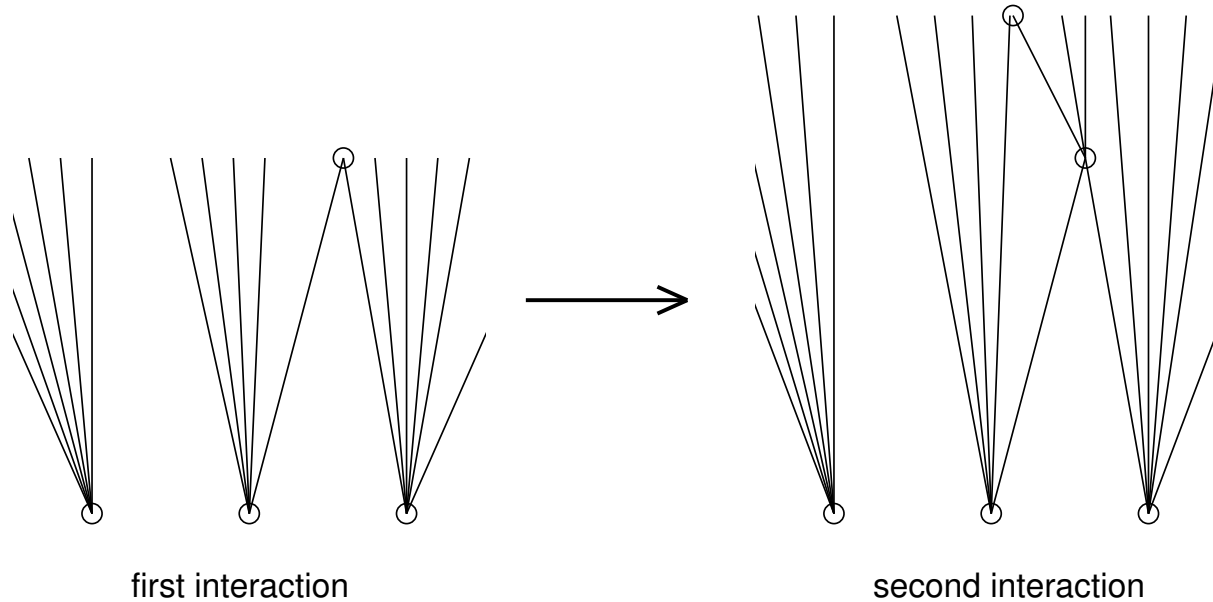
Example 1 cont'd



Example 1 cont'd



The Cauchy Problem: Front Tracking



Start ($t = 0$): piecewise constant initial data

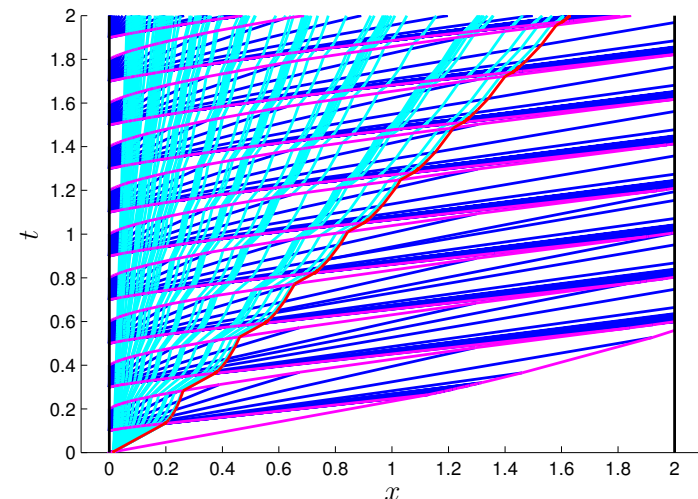
→ sequence of local Riemann problems

→ p.w discontinuities between (x,t) -rays

While $t < t_{end}$:

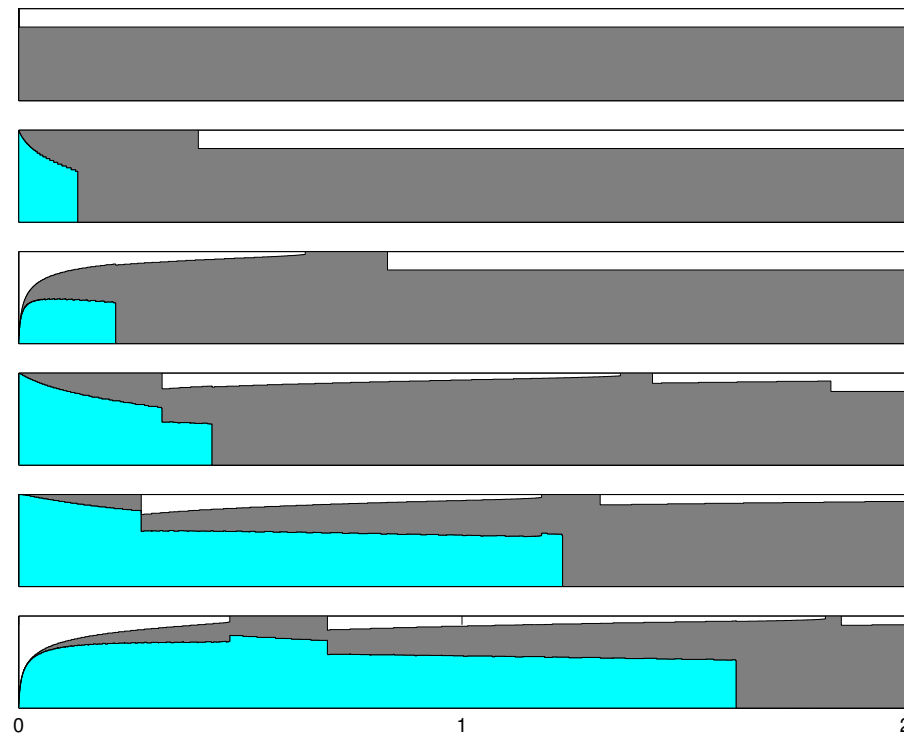
track discontinuities

solve Riemann problems



Example 2

- Initially, reservoir filled with 80% oil and 20% gas
- Alternate cycles of water and gas injection
- Front-tracking solution (with $\Delta_u = 0.005$)
- Half a million Riemann solves ~ 5 sec on a desktop PC



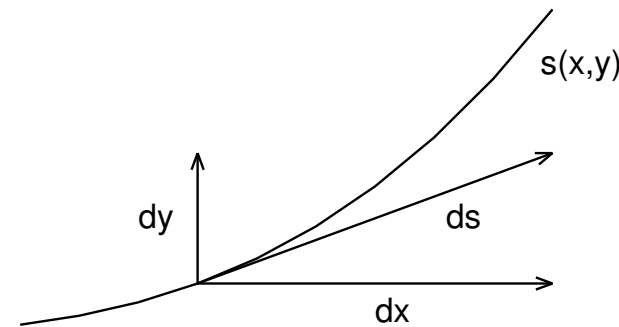
Streamline Methods

Interpret the saturation equation $\phi \partial_t S + v \cdot \nabla f(S) = 0$ as an equation along streamlines using

$$\frac{v}{|v|} = \left[\frac{dx}{ds}, \frac{dy}{ds}, \frac{dz}{ds} \right]^T \quad \text{or} \quad v \cdot \nabla = |v| \frac{\partial}{\partial s}$$

Transformation using *time-of-flight* τ

$$|v| \frac{\partial}{\partial s} = \phi \frac{\partial}{\partial \tau}$$



gives a family of 1-D transport equations along streamlines

$$\partial_t S + \partial_\tau f(S) = 0.$$

Streamline Simulation

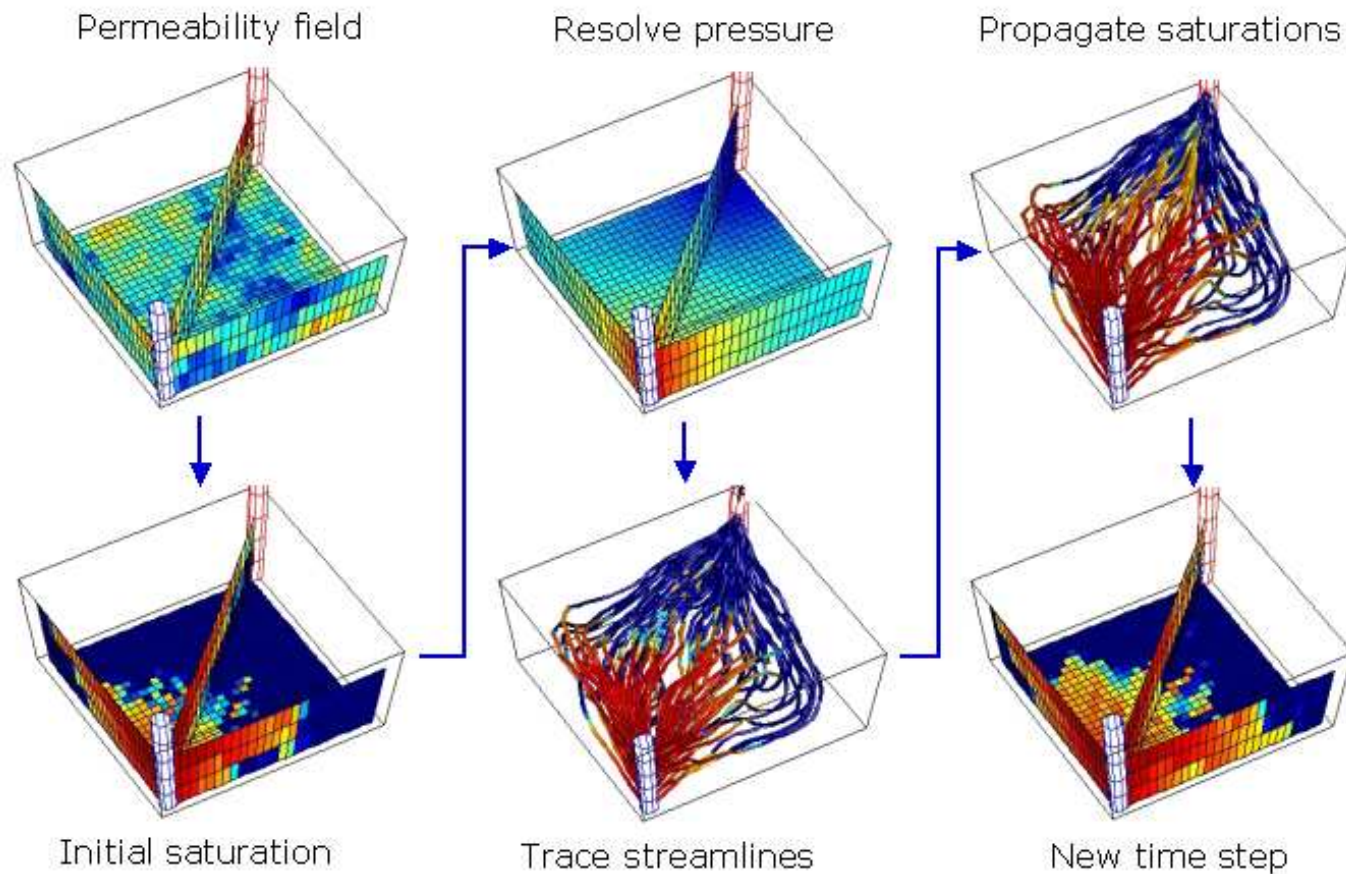
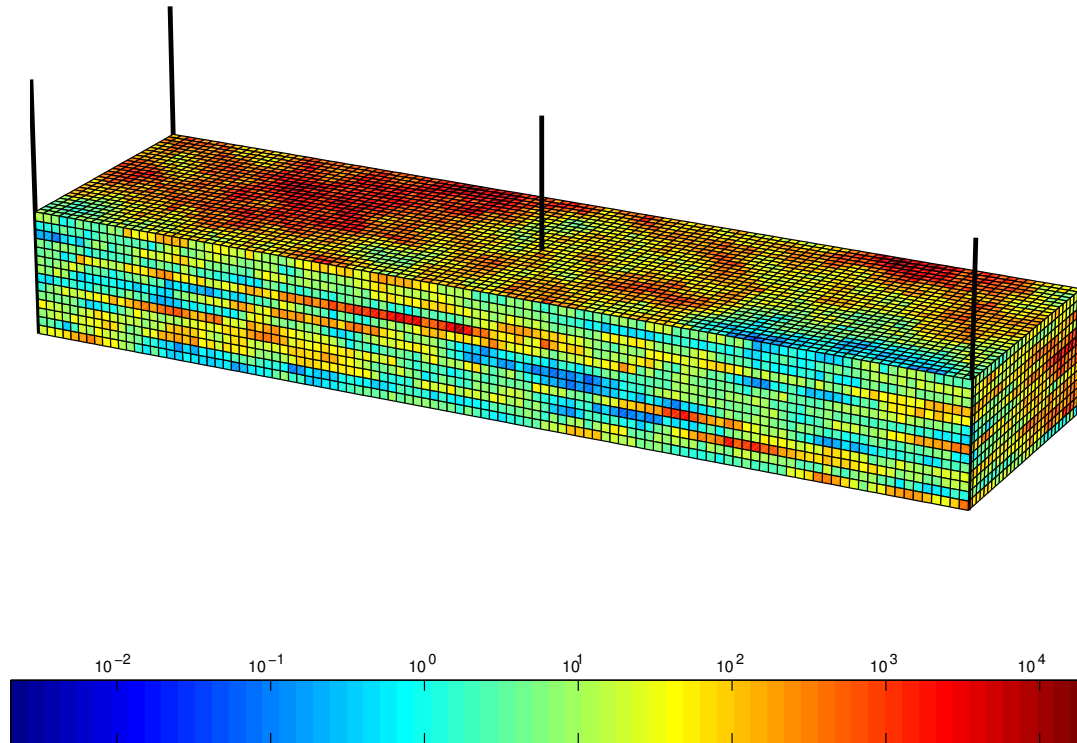


Figure from Yann Gautier

Example 4: SPE 10 Tarbert Formation

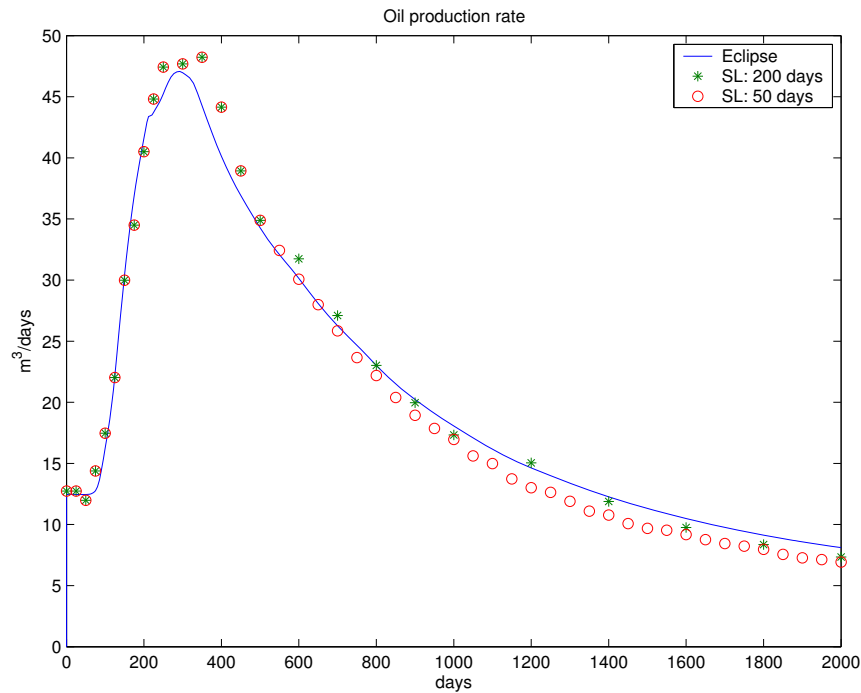


- $30 \times 110 \times 15$ upscaled sample from Tarbert formation
- Initial composition: $(S_w, S_g) = (0.0, 0.2)$
- 2000 days of production
- Either: continuous water injection
- Or: water-alternating-gas every 200 day

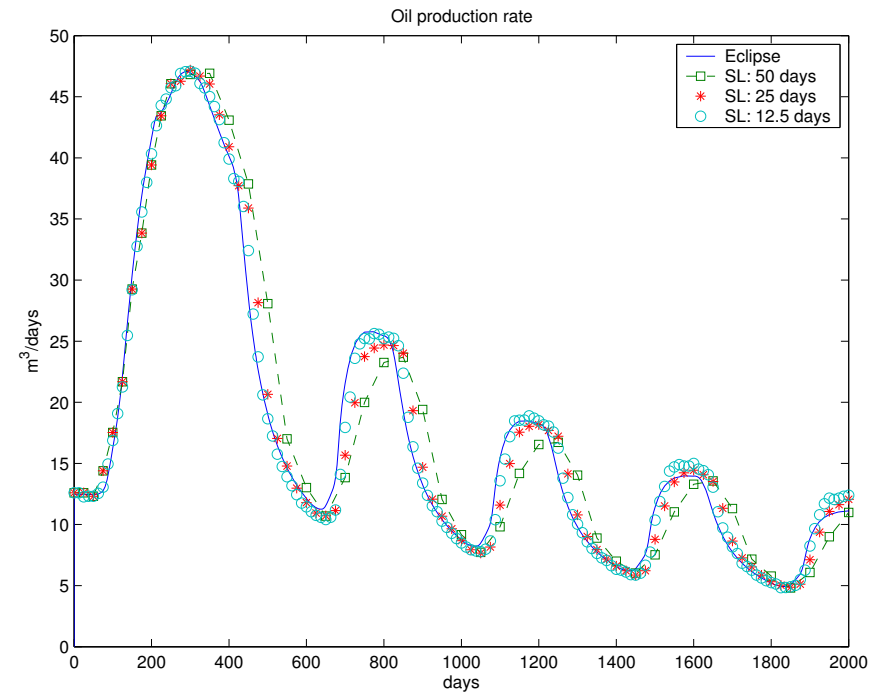
Data reduction is used to speed up front-tracking solution: weak interactions are treated as being of type $\mathcal{S}_1 \mathcal{S}_2$.

Example 4 cont'd

Oil production:



water injection



water alternating gas injection

Runtimes: 8 hr 20 min for Eclipse, 2 hr 13 min for streamlines