

Multiscale Mixed Finite Elements for the Stokes–Brinkman Equations

K.-A. Lie and A. F. Gulbransen

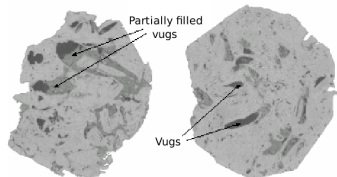
SINTEF ICT, Dept. Applied Mathematics, Norway

ENUMATH 2009, Uppsala, Sweden

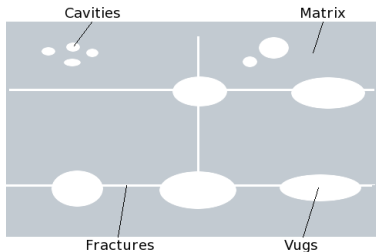
Motivation: Carbonate Reservoirs

Carbonate reservoirs contain:

- 60% of the world's oil reserves
- 40% of the world's gas reserves
- consist of free-flow and porous regions



(Liying Zhang, 2005)



(courtesy of NTNU)

Conventional Approach:

Porous region:

Darcy's law, mass conservation:

$$\begin{aligned}\mu \mathbf{K}^{-1} \vec{u}_D + \nabla p_D &= \vec{f} && \text{in } \Omega_D \\ \nabla \cdot \vec{u}_D &= q && \text{in } \Omega_D\end{aligned}$$

Interface conditions ⇔

Free-flow region:

Stokes equations:

$$\begin{aligned}-\mu \Delta \vec{u}_S + \nabla p_S &= \vec{f} && \text{in } \Omega_S \\ \nabla \cdot \vec{u}_S &= q && \text{in } \Omega_S\end{aligned}$$

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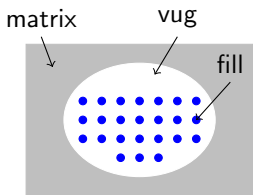
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Problems:

- domains not well separated
- difficulties obtaining precise information about location and geometry
- hard to model loose fill-in material

The Stokes–Brinkman model

Introduce a single-parameter family

$$\begin{aligned}\mu \mathbf{K}^{-1} \vec{u} + \nabla p - \tilde{\mu} \Delta \vec{u} &= \vec{f} & \text{in } \Omega \\ \nabla \cdot \vec{u}_S &= q & \text{in } \Omega.\end{aligned}$$

$\tilde{\mu}$ – effective viscosity, μ – fluid viscosity

Special cases:

- $\mathbf{K} \rightarrow \infty, \tilde{\mu} = \mu \implies$ Stokes–Brinkman \longrightarrow Stokes
- $\tilde{\mu} = 0 \implies$ Stokes–Brinkman \longrightarrow Darcy

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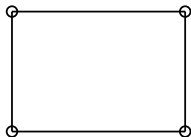
- $\mathbf{K} \rightarrow \infty, \tilde{\mu} = \mu \implies$ Stokes–Brinkman \longrightarrow Stokes
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Here: $\tilde{\mu} = \mu$. For typical parameters seen in carbonate reservoirs

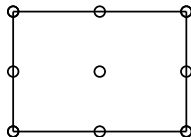
$$\nabla p = -\mu \mathbf{K}^{-1} \vec{u} + \tilde{\mu} \Delta \vec{u} \approx -\mu \mathbf{K}^{-1} \vec{u}$$

2-dimensional Taylor-Hood elements:

Pressure (Q_1):

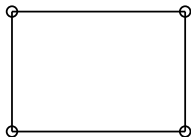


Velocity (Q_2):

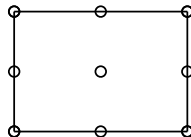


2-dimensional Taylor-Hood elements:

Pressure (Q_1):



Velocity (Q_2):



Variational formulation:

Find $\vec{u} \in V$ and $p \in Q$ such that

$$b(\vec{u}, \vec{v}) - c(p, \vec{v}) = 0 \quad \forall \vec{v} \in V$$

$$c(\vec{u}, \pi) = (q, \pi) \quad \forall \pi \in Q$$

where $V \subset (H^1(\Omega))^2$ and $Q \subset L^2(\Omega)$

Mixed finite-element system (Stokes-Brinkman)

$$\begin{bmatrix} \mathbf{B}_1 & \mathbf{0} & \mathbf{C}_1 \\ \mathbf{0} & \mathbf{B}_2 & \mathbf{C}_2 \\ \mathbf{C}_1^\top & \mathbf{C}_2^\top & \mathbf{0} \end{bmatrix} \begin{bmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \\ -p \end{bmatrix} = \begin{bmatrix} \mathbf{0} \\ \mathbf{0} \\ q \end{bmatrix}$$

The entries in the matrices are:

$$B_{ij,k} = \int_{\Omega} \mu v_i K_k^{-1} v_j \, d\Omega + \int_{\Omega} \tilde{\mu} \left(\frac{\partial v_i}{\partial x_1} \frac{\partial v_j}{\partial x_1} + \frac{\partial v_i}{\partial x_2} \frac{\partial v_j}{\partial x_2} \right) \, d\Omega,$$

$$C_{ij,k} = \int_{\Omega} \frac{\partial v_i}{\partial x_k} \pi_j \, d\Omega.$$

where $k = 1, 2$ denotes the spatial dimension and $\mathbf{K} = \begin{bmatrix} \mathbf{K}_1 & \mathbf{0} \\ \mathbf{0} & \mathbf{K}_2 \end{bmatrix}$.

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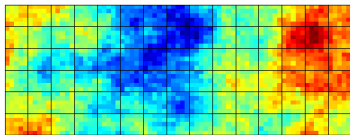
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100 × 100 cells ⇒ 91.003 dofs ⇒ multiscale multiphysics method

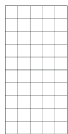
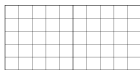
Feasible Approaches: Upscaling or Multiscale Methods

Darcy's law on coarse scale, Stokes–Brinkman on fine scale

Flow-based upscaling:



Coarse blocks (Darcy):



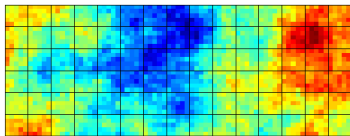
Darcy–Stokes or
Stokes–Brinkman:



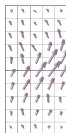
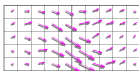
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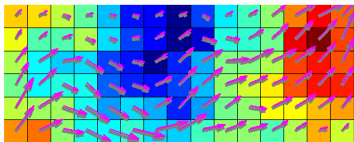
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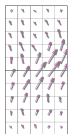
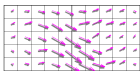
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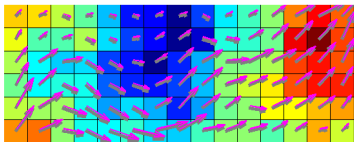
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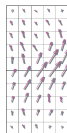
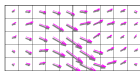
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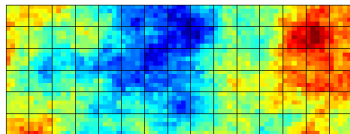
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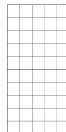
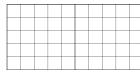
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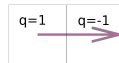
Multiscale method:



Coarse blocks (Darcy):



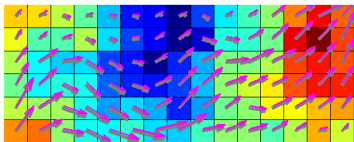
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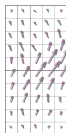
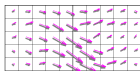
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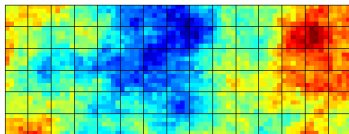
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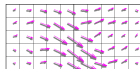
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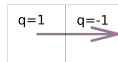
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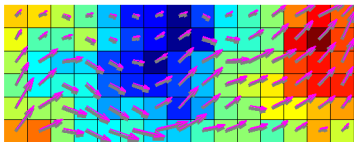
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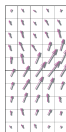
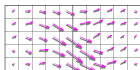
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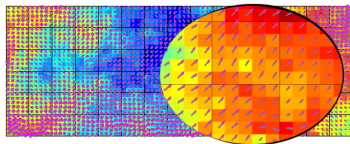
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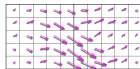
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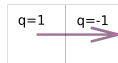
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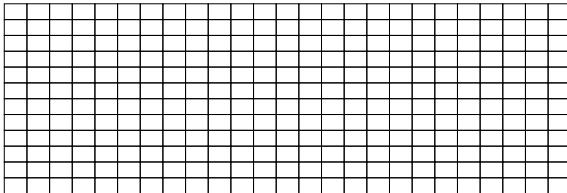


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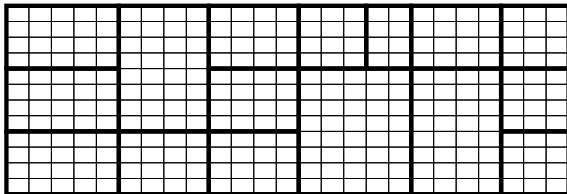
Grids and basis functions:

Fine grid with permeability attached to each cell:



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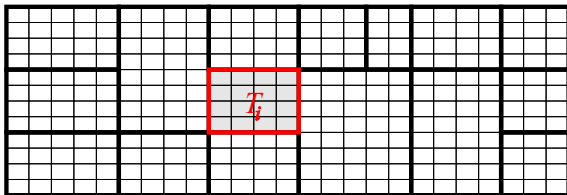
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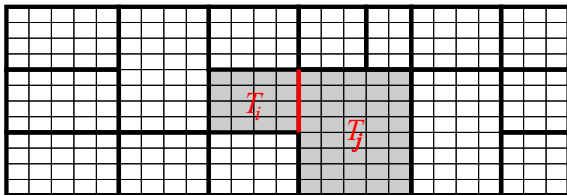


Construct a *coarse* grid, and choose the discretisation spaces U and V^{ms} such that:

- For each coarse block T_i , there is a basis function $\phi_i \in U$.

Grids and basis functions:

Fine grid with permeability attached to each cell:



Construct a *coarse* grid, and choose the discretisation spaces U and V^{ms} such that:

- For each coarse block T_i , there is a basis function $\phi_i \in U$.
- For each coarse edge Γ_{ij} , there is a basis function $\psi_{ij} \in V^{\text{ms}}$.

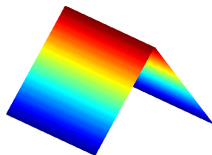
Decomposition:

- $p(x, y) = \sum_i p_i \phi_i(x, y)$ – sum over all coarse blocks
- $v(x, y) = \sum_{ij} v_{ij} \psi_{ij}(x, y)$ – sum over all block faces

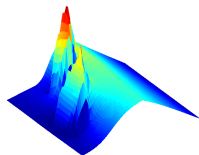
Basis ϕ_i for pressure:

$$\phi_i = \begin{cases} 1 & \text{in } T_i, \\ 0 & \text{otherwise.} \end{cases}$$

Basis ψ_{ij} for velocity:



homogeneous (RT0)



heterogeneous

Construction of Multiscale Basis Functions

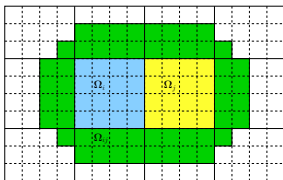
Velocity basis function ψ_{ij} solves a local system of equations:

$$\mu \mathbf{K}^{-1} \vec{\psi}_{ij} + \nabla \varphi_{ij} - \tilde{\mu} \Delta \vec{\psi}_{ij} = 0,$$

$$\nabla \cdot \vec{\psi}_{ij} = \begin{cases} w_i(\vec{x}), & \text{if } \vec{x} \in T_i, \\ -w_j(\vec{x}), & \text{if } \vec{x} \in T_j, \\ 0, & \text{otherwise.} \end{cases}$$

$w_i \propto \text{trace}(K_i)$ drives a unit flow through Γ_{ij} .

If there is a sink/source in T_i , then $w_i \propto q_i$.



Construction of Multiscale Basis Functions

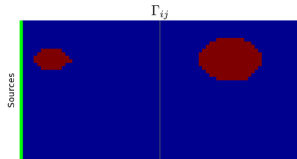
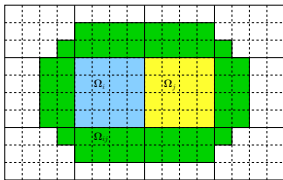
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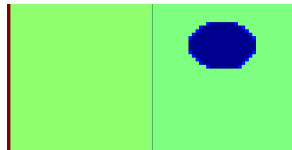
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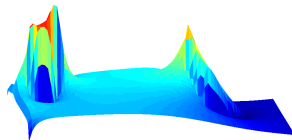
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Permeability



Weight function

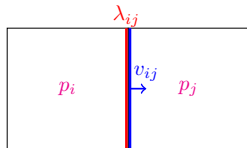


Basis function

Hybrid Formulation (for the Darcy Equations)

Mixed method gives a saddle-point problem \longrightarrow hybrid formulation:

$$\begin{bmatrix} B & C & D \\ C^T & 0 & 0 \\ D^T & 0 & 0 \end{bmatrix} \begin{bmatrix} v \\ -p \\ \lambda \end{bmatrix} = \begin{bmatrix} 0 \\ f \\ 0 \end{bmatrix}$$



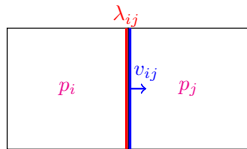
where:

$$B_{ij} = \int_{\Omega} \mu \vec{v}_i \mathbf{K}^{-1} \vec{v}_j d\Omega, \quad C_{ij} = \int_{\Omega} \chi_{T_j} \nabla \cdot \vec{v}_i d\Omega, \quad D_{ij} = \int_{\partial\Omega} |\vec{v}_i \cdot \vec{n}_j| ds,$$

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Split the basis functions, $\psi_{ij} = \psi_{ij}^H - \psi_{ji}^H$

$$\psi_{ij}^H(E) = \begin{cases} \psi_{ij}(E), & \text{if } E \in T_{ij} \setminus T_j \\ 0, & \text{otherwise} \end{cases} \quad \psi_{ji}^H(E) = \begin{cases} -\psi_{ij}(E), & \text{if } E \in T_j \\ 0, & \text{otherwise} \end{cases}$$

Hybrid basis functions ψ_{ij}^H as columns in a matrix Ψ

MsMFEs for the Stokes–Brinkman Equations

Coarse-scale hybrid mixed system (Darcy):

$$\begin{bmatrix} \hat{\Psi}^T B^{TH} \hat{\Psi} & \hat{\Psi}^T C \mathcal{I} & \hat{\Psi}^T D \mathcal{J} \\ \mathcal{I}^T C^T \hat{\Psi} & \mathbf{0} & \mathbf{0} \\ \mathcal{J}^T D^T \hat{\Psi} & \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{bmatrix} v^c \\ -p^c \\ \lambda^c \end{bmatrix} = \begin{bmatrix} \mathbf{0} \\ f^c \\ \mathbf{0} \end{bmatrix}$$

$$\hat{\Psi} = \Psi A^{-1}$$

Ψ – matrix with basis functions

A – matrix with face areas

B^{TH} – fine-scale Darcy TH-discretization

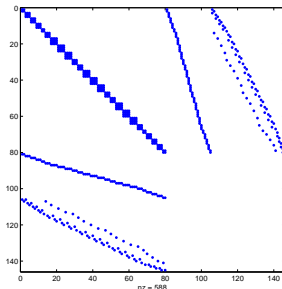
\mathcal{I} – prolongation from blocks to cells

\mathcal{J} – prolongation from block faces to cell faces

Reconstruction of fine-scale velocity

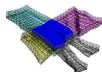
$$v^f = \hat{\Psi} v^c$$

(Pressure bases may also have fine-scale structure if necessary)

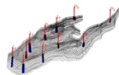


Test of the multiscale Darcy/Stokes-Brinkman method:

- 1 2D sandstone reservoirs (no free-flow regions)
- 2 2D vuggy reservoir (short correlation)
- 3 2D fractured reservoir (long correlation)
- 4 2D vuggy and fractured reservoir (short and long correlation)
- 5 3D core sample

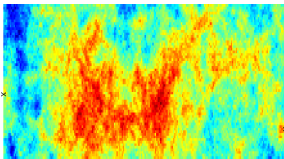


All simulations in the MATLAB Reservoir Simulation Toolbox
<http://www.sintef.no/MRST> (GNU Public License)

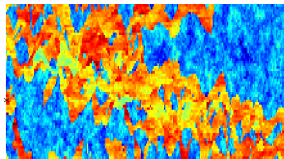


Example 1: Sandstone reservoir

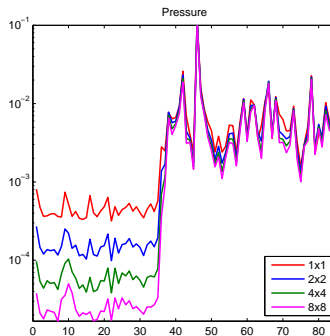
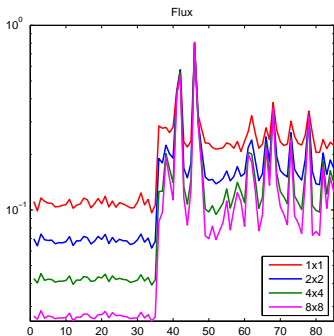
Model 2 of the 10th SPE Comparative Solution Project



Tarbert (1-35)

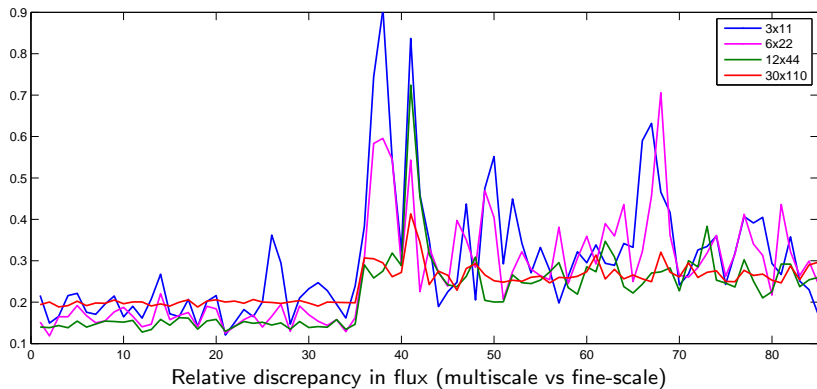


Upper Ness (36-85)

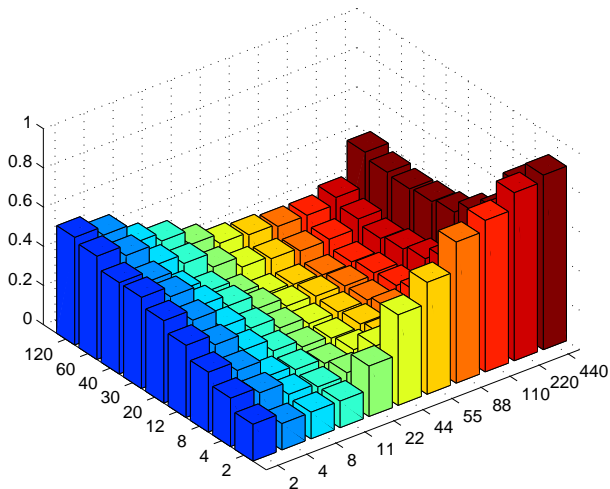


Relative discrepancy between RT0 and Q_2/Q_1

Example 1: Sandstone reservoir

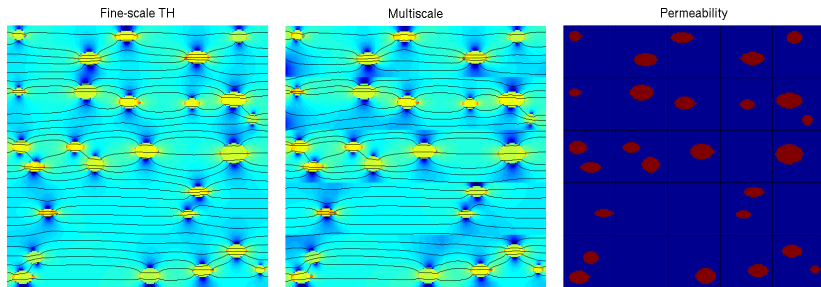


Example 1: Sandstone reservoir



Relative discrepancy in flux (multiscale vs fine-scale)
 2×2 refinement of Layer 21

Example 2: Vuggy reservoir (short correlation)



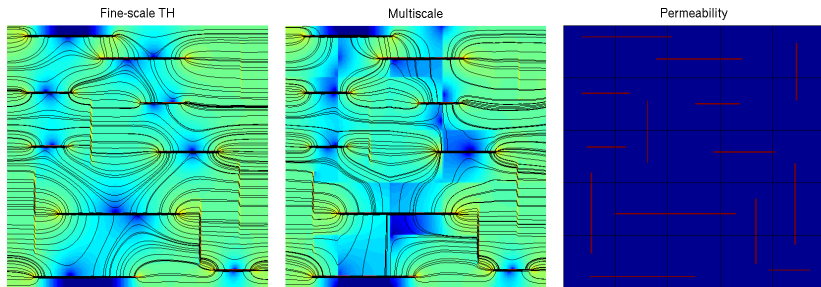
Fine-scale model:
 200×200 cells

Multiscale model:
 5×5 blocks

$$K_{vugs} = 10^7 \times K_{matrix}$$

26 random vugs (areas = 1.8–10.4 m²), $\frac{\|v^{ms} - v\|_2}{\|v\|_2} = 0.07$

Example 3: Fractured reservoir (long correlation)



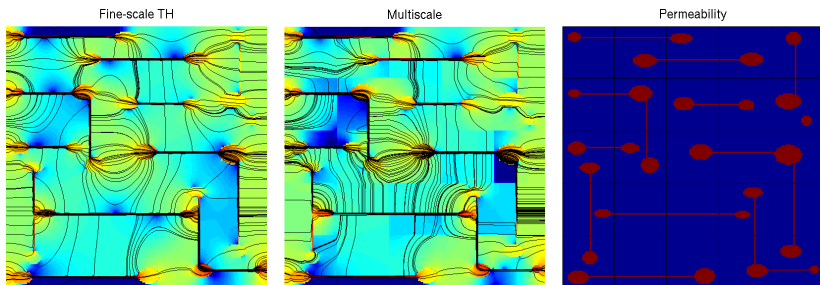
Fine-scale model:
 200×200 cells

Multiscale model:
 5×5 blocks

$$K_{vugs} = 10^7 \times K_{matrix}$$

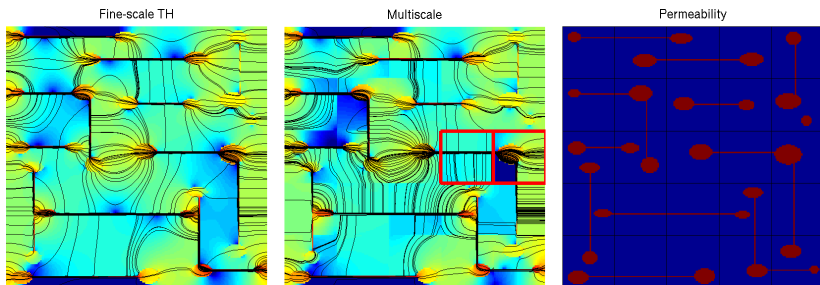
14 random fractures of varying length, $\frac{\|v^{ms} - v\|_2}{\|v\|_2} = 0.07$

Example 4: Vuggy and fractured reservoir



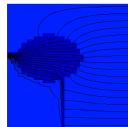
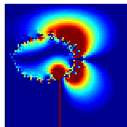
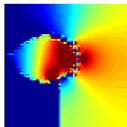
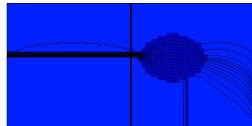
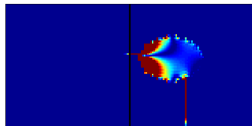
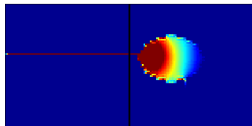
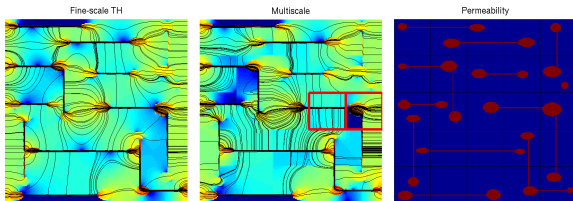
$$\frac{\|v^{ms} - v\|_2}{\|v\|_2} = 0.09$$

Example 4: Vuggy and fractured reservoir



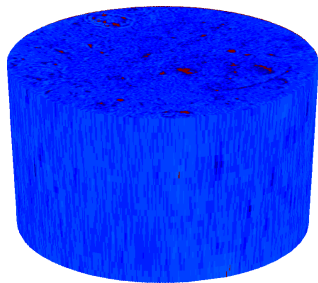
$$\frac{\|v^{ms} - v\|_2}{\|v\|_2} = 0.09$$

Example 4: Vuggy and fractured reservoir



Example 5: Core sample from Shell E&P

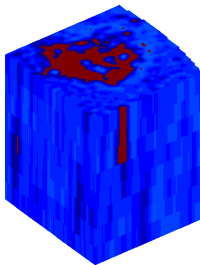
Full model:



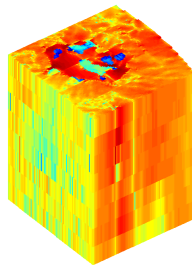
$512 \times 512 \times 26$ cells
3.449.654 active

Subsample:

Log10 permeability [Darcy]



Log10 MS flux [m/day]



$85 \times 85 \times 8$ cells, 55.192 active, 75 blocks
pressure boundary conditions

Proof of concept for our MsMFE method:

- Multiphysics with different equations on the coarse and fine scales
- Can be used to simulate flow in carbonate reservoirs

Challenges/issues:

- Very high number of dofs
- Raviart–Thomas not exactly reproduced for homogeneous medium

Ideas to pursue:

- Triangular and unstructured grids
- Correction functions
- Use of (limited) global information